

# *Avenue to spin foam phenomenology* – **Generalized spacetime geometries and continuum effective dynamics**

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• spin foams = covariant path integrals for quantum gravity

 $\mathcal{D}\mu(\text{geometry})\exp(iS[\text{geometry}])$ 

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geometry classical: [g]
geometry quantum: fuzzy
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• phenomenon relevant for large class of spin-foam models:

(primary) simplicity constraints classical: {1st class} (primary) simplicity constraints quantum: {1st class}  $\cup$  {2nd class, {  $\cdot$  ,  $\cdot$  }  $\sim \gamma$ } — imposed weakly

• consequence: path integral taken over extended configuration space

geometry semi-classical (effective): ?

#### <u>Hypothesis</u>: Semi-classical ( $\leftrightarrow$ extended classical) configuration space is captured by area metrics.

#### plan:

- area metric
- evidence for area metrics in semi-classical limit of spin foams
- area-metric dynamics from "classical constraints imposed weakly"

### **Motivation**

- effective length-metric action
- geometric interpretation of non-metric d.o.f.
- $\Rightarrow$  prospect for phenomenology & ...





### **Area metrics**

**area metric** [Schuller & Wohlfahrt, 2005; Punzi, Schuller & Wohlfarth, 2006]

$$G_p: \Lambda^2 T_p \mathcal{M} \times \Lambda^2 T_p \mathcal{M} \to \mathbb{R} \quad \in \mathcal{T}_p^{(0,4)}(\mathcal{M})$$

i) 
$$G_{\mu\nu\rho\sigma} = G_{\rho\sigma\mu\nu}$$
  
ii)  $G_{\mu\nu\rho\sigma} = -G_{\nu\mu\rho\sigma} \ (= -G_{\mu\nu\sigma\rho})$   
iii)  $G_{IJ}, \ I, J = [\mu\nu] = 1, \dots, d(d-1)/2$  non-degenerate

 $G(A, A) = ||A||^2$  area of parallelogram A (viewed as simple bivector  $A = l \wedge r$ )  $G(l \land t, r \land t) \leftrightarrow 3d$  dihedral angle between planes ( $\bot$  bivectors) intersecting at a line

cyclic area metrics

iv) 
$$G_{\mu[\nu\rho\sigma]} = 0$$

d = 4: 21 - 1 = 20 d.o.f.



length metric  $g_p: T_p\mathcal{M} \times T_p\mathcal{M} \to \mathbb{R} \quad \in \mathcal{T}_p^{(0,2)}(\mathcal{M})$ i)  $g_{\mu\nu} = g_{\nu\mu}$ ii)  $g_{\mu\nu}$ ,  $\mu, \nu = 1, ..., d$  non-degenerate  $g(l, l) = ||l||^2$  length of vector l

 $g(r, l) \leftrightarrow 2d$  angle between vectors l, r intersecting at a point

induced area metric

$$G^g_{\mu\nu\rho\sigma} = g_{\mu\rho}g_{\nu\sigma} - g_{\mu\sigma}g_{\nu\rho}$$

d = 4 : 10 d.o.f.

 $g_{eff}$ ?





### **Area Regge action — continuum limit**

 $l_e$ : length of edge e $A_t$ : area of triangle t $\theta_{\sigma,t}$ : dihedral angle at triangle *t* in simplex  $\sigma$  $\epsilon_t = 2\pi k - \sum \theta_{\sigma,t}$ : deficit angle at triangle *t* 

length Regge calculus [Regge, 1961]

 $S_{LR}[l_e] = \sum A_t(l_e)\epsilon_t(l_e) \iff$ **vacuum GR** on piecewise flat  $\Delta^{(d=4)}$ ,  $\epsilon_t \leftrightarrow$ curvature change of variables  $l_e \rightarrow A_t$  invertible only locally

area Regge calculus [Rovelli, 1993; Barret, Rocek & Williams, 1999; Asante, Dittrich & Haggard, 2018]

 $S_{AR}[A_t] = \sum A_t \epsilon_t(A_t) \quad \Leftrightarrow \quad \text{dynamics of many more d.o.f.}, \quad \epsilon_t \leftrightarrow \text{curvature} + \text{shape mismatch}$  $\epsilon_t(A_t) = 0$  e.o.m. do not imply flatness! describes semi-classical limit of spin foams

lattice continuum limit of Area-Regge action: [Dittrich, 2021; Dittrich & Kogios, 2022]

 $S_{AR}[A_t] \xrightarrow{\min_{\lambda \to 0}} S[G]$  dynamics of area metric *G* associated to hypercube,  $G \leftrightarrow g$  (massless graviton) + massive d.o.f.

$$S_{eff}[g] = S_{EH}[g] + Weyl^2 + \mathcal{O}(\lambda^4)$$





 $n_{l_e} = n_{A_t} = 10$ 

one simplex:



glued simplices:  $n_{A_t} > n_{l_e}$ 



# **Semi-classical simplex geometry**

 $\tau$ : tetrahedron  $\sigma$ : 4-simplex

[Dittrich & Padua-Argüelles, 2023]

microscopic identification: semi-classical spin foam d.o.f.  $\leftrightarrow$  area-metric G d.o.f.

- classical  $\sigma$ : 10 d.o.f. = 10 x lengths or 10 x areas
- quantum  $\sigma$ : 5 quantum  $\tau$  glued as dictated by quantum uncertainty quantum  $\tau$ : 4 x norms of normal vectors (4 x areas) + 1 x inner product between two normals (1 x 3d dihedral angle) only areas of the 10 triangles shared by pairs of tetrahedra can be identified  $15 \text{ d.o.f.} = 5 \times 5 - 10$
- **semi-classical/coherent/twisted**  $\sigma$ : 5 classical  $\tau$  with quantum glue classical  $\tau$ : 6 x lengths or 4 x areas + 2 x 3d dihedral angles at non-opposite edges only areas of the 10 triangles shared by pairs of tetrahedra are identified 20 d.o.f. = 5 x 6 – 10  $\leftrightarrow$  10 x areas + 10 x 3d dihedral angles (2 x 3d dihedral angles per  $\tau$ , at non-opposite edges)

**20 d.o.f. of a semi-classical**  $\sigma$  can be arranged into 20 d.o.f. for an area metric *G* associated to  $\sigma$  [Dittrich & Padua-Argüelles, 2023]





# $\Rightarrow$ Microscopic & macroscopic area metrics appear in the semi-classical regime of spin foams.

### Goal:

constraints.

- Parametrize the continuum classical extended configuration space by area metrics, starting from the configuration variables in the Plebanski formulation of GR ( $\leftrightarrow$  origin of spin-foam quantization).
- Reproduce the continuum effective dynamics of spin foams by weakening a subset of the classical





### **Plebanski formulation of general relativity**

### non-chiral Plebanski action [Plebanski, 1977; Capovilla, Jacobson, Dell & Mason, 1991, Reisenberger, 1995; Pietri & Freidel, 1999]

*B* field: so(4) valued 2-form  $\omega$  connection: so(4) valued 1-form  $\phi$ : Lagrange multiplier *γ*: Barbero-Immirzi parameter

$$S_{P}[B,\omega,\phi] = \int \underbrace{\delta_{IJKL}B^{IJ} \wedge F^{KL}(\omega)}_{\mathbf{BF}} \underbrace{+ \frac{1}{2\gamma} \epsilon_{IJKL}B^{IJ} \wedge F^{KL}(\omega)}_{\mathbf{Holst}} \underbrace{- \frac{1}{2} \phi_{IJKL}B^{IJ} \wedge B^{KL}}_{\delta\phi: \text{ constraints}}, \qquad \phi_{IJKL} = \phi_{KLIJ} = -\phi_{JIKL}, \ \phi_{IJKL} \epsilon^{IJKL} = 0$$

#### simplicity constraints

$$B^{IJ} \wedge B^{KL} = \nu \epsilon^{IJKL}, \quad \nu = \frac{1}{4!} B^{AB} \wedge B^{CD} \epsilon_{ABCD} \quad (\nu \neq 0)$$

4 solution sectors

 $B^{IJ} = \pm (\star e \wedge e)^{IJ}$  gravitational sector (—Palatini action for GR)  $B^{IJ} = \pm (e \wedge e)^{IJ}$  topological sector

. constraints υψ



### **Modified Plebanski theories**

**Plebanski action** [Plebanski, 1977; Capovilla, Jacobson, Dell & Mason, 1991] BF (+ Holst) + simplicity constraints ( $\phi$  Lagrange multiplier)

modified Plebanski action [Krasnov, 2006-9; Bengtsson, 2007] BF (+ Holst) + potential  $V(\phi)$ 

e.o.m. from varying  $\phi: B \wedge B = \nu \left( \epsilon + \frac{\partial V}{\partial \phi} \right) \xrightarrow{\det(V^{(2)}) \neq 0} \phi(B)$   $S_{P, mod.}[B, \omega] = \left[ \delta_{IJKL} B^{IJ} \wedge F^{KL}(\omega) + \frac{1}{2\nu} \epsilon_{IJKL} B^{IJ} \wedge F^{KL}(\omega) + V(B), \quad \omega(B) \text{ solution to } d_{\omega}B = 0 \longrightarrow S_{P, mod. eff.}[B] \text{ second-order} \right]$ [Freidel, 2008; Speziale, 2010; Beke, Palmisano & Speziale, 2012]

non-degenerate case det  $(V^{(2)}) \neq 0 \Leftrightarrow$ full set of simplicity constraints  $\mapsto$  potential for *B* chiral (self-dual),  $\mathscr{L}(G) = \operatorname{su}(2)$ : 2 massless propagating d.o.f. = graviton, "deformations of GR" [Krasnov, 2006-9; Freidel, 2008] non-chiral,  $\mathscr{L}(G) = so(4)$ : 8 = 2 + 6 d.o.f. [Alexandrov & Krasnov, 2008]

degenerate case det  $(V^{(2)}) = 0 \Leftrightarrow$ full set of simplicity constraints  $\mapsto$  subset of simplicity constraints (sharp) + potential for rest of *B* (suppressed)

**goal:** use mechanism to derive area-metric action from modified non-chiral Plebanski action:  $S_{P, mod. eff.}[B_{\mu\nu}^{IJ}] \rightarrow S[G]$  [JB & Dittrich, 2022]





#### counting d.o.f.

 $B_{\mu\nu}^{IJ}$ : 36 = 30 + 6 (SO(4) gauge) d.o.f.  $\delta \phi$ : 20 simplicity constraints (2,0)  $\oplus$  (0,2)  $\oplus$  (1,1)  $\oplus$  (0,0)

 $e_{\mu}^{I}$ : 16 = 10 + 6 (SO(4) gauge) d.o.f.

 $g_{\mu\nu}$  : 10 d.o.f.

#### parametrization of B-field

i)  $so(4) = su(2)_{+} \oplus su(2)_{-}$  isomorphism,  $P_{\pm i}^{IJ} = \pm \delta_{0i}^{IJ} + \frac{1}{2}\epsilon_{0i}^{IJ}$  projectors:  $B_{\mu\nu}^{IJ} = P_{+i}^{IJ}B_{+\mu\nu}^{i} + P_{-i}^{IJ}B_{-\mu\nu}^{i}$ 

ii) parametrization of su(2) B-field  $B_+$  using Urbantke theorem: 4d metric  $g_+$  & unimodular 3x3 matrix of scalars  $b_+$ [Urbantke, 1984; Freidel, 2008]

$$B^i_{\pm\,\mu\nu} = b^i_{\pm\,a} \Sigma^a_{\pm\,\mu\nu}(e_{\pm})$$

 $\Sigma^{a}_{\pm\mu\nu}(e_{\pm}) = \pm e^{0}_{\pm[\mu}e^{i}_{\pm\nu]} + \epsilon^{i}_{jk}e^{j}_{\pm\mu}e^{k}_{\pm\nu}$  Plebanski 2-forms (self-/antiself-dual) — encode 2 x spacetime metrics:  $g_{\pm\mu\nu} = e^{I}_{\pm\mu}e^{J}_{\pm\nu}\delta_{IJ}$  $b_{\pm}$  define 2 x unimodular internal metrics:  $q_{\pm ab} = b_{\pm a}^{i} b_{\pm b}^{j} \delta_{ij}$ 

 $B_{\mu\nu}^{IJ}(36) \Leftrightarrow g_{+\mu\nu}(10), g_{-\mu\nu}(10), b_{+a}^{i}(8), b_{-a}^{i}(8)$ 

 $S_{P, mod. eff.}[B_{\mu\nu}^{IJ}] = S_{P, mod. eff.}[e_{+\mu}^{I}, e_{-\mu}^{I}, q_{+ab}, q_{-ab}] = S_{BF eff., su(2)_{+}}[e_{+\mu}^{I}, q_{+ab}, q_{-ab}]$ 

 $B_{\mu\nu}^{IJ}$ : 36 = 30 + 6 (SO(4) gauge) d.o.f.  $\delta \phi'$ : 10 simplicity constraints — which subset?

 $G_{\mu\nu\rho\sigma}$  : 20 (gauge-invariant) d.o.f.

$$[e_{+\mu}^{I}, q_{+ab}] + S_{BF\,eff.,\,su(2)}[e_{-\mu}^{I}, q_{-ab}] + \int V[e_{+\mu}^{I}, e_{-\mu}^{I}, q_{+ab}, q_{-ab}] \quad \{\text{Spez}\}$$



$$B_{\mu\nu}^{IJ}(36) \quad \Leftrightarrow \quad g_{+\mu\nu}(10), g_{-\mu\nu}(10), b_{+a}^{i}(8), b_{-a}^{i}(8)$$

#### area metric from B-field

define cyclic left- and right-handed area metrics:  $G_{\pm\,\mu\nu\rho\sigma} \stackrel{\cdot}{\equiv} \frac{1}{2} \left( (B^{i}_{\pm} \otimes B_{\pm\,i})_{\mu\nu\rho\sigma} - \frac{1}{\Delta\,!} (B^{i}_{\pm} \otimes B_{\pm\,i})_{\alpha\beta\gamma\delta} \epsilon^{\alpha\beta\gamma\delta} \epsilon_{\mu\nu\rho\sigma} \right)$ 

imposing equal left- and right-handed area metrics: [Reisenberge  $G_{+\mu\nu\rho\sigma} \stackrel{!}{=} G_{-\mu\nu\rho\sigma} \Leftrightarrow 20 \text{ simplicity constraints } (0,0) \oplus (1,1)$ 

imposing flat left- and right-handed internal metrics: [Speziale, 2010; Beke, Palmisano & Speziale, 2012]  $q_{\pm ab} \stackrel{!}{=} \delta_{ab} \quad \Leftrightarrow \quad 10 \text{ simplicity constraints } (0,0) \oplus (1,1) \oplus (2,0) \oplus (0,2) \quad \longrightarrow \quad S[g_{+\mu\nu},g_{-\mu\nu}], \text{ bi-metric gravity (unstable)}$ 

imposing equal left- and right-handed length metrics: [JB & Dittrich, 2022]  

$$g_{+\mu\nu} \stackrel{!}{=} g_{-\mu\nu} \Leftrightarrow \mathbf{10}$$
 simplicity constraints  $(0,0) \oplus (1,1) \oplus (2,0) \oplus (0,2) \longrightarrow S[g_{\mu\nu}, q_{+ab}, q_{-ab}] \triangleq S[G_{\mu\nu\rho\sigma}]$ , area-metric actives  
 $G_{\mu\nu\rho\sigma} \equiv G_{+\mu\nu\rho\sigma} + G_{-\mu\nu\rho\sigma} = \frac{1}{2} \sum_{\pm} q_{\pm ab} \Sigma^{a}_{\mu\nu}(e) \Sigma^{b}_{\rho\sigma}(e) - \frac{1}{6} \operatorname{Tr}(q_{+} - q_{-}) \det(e) \epsilon_{\mu\nu\rho\sigma}$   
 $G_{\mu\nu\rho\sigma}(20) \Leftrightarrow g_{\mu\nu}(10), q_{+ab}(5), q_{-ab}(5)$   
 $q_{\pm ab}$  scalars  
 $g_{Urbantke}$ 

(er, 1995]  
() 
$$\oplus$$
 (2,0)  $\oplus$  (0,2)  $\longrightarrow$   $S[g_{\mu\nu}]$ , **GR**



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$$S[G_{\mu\nu\rho\sigma}] \stackrel{\circ}{=} S_{P, mod. eff.}[e^{I}_{\mu}, q_{+ab}, q_{-ab}] = S_{BF \, eff., \, su(2)_{+}}[e^{I}_{\mu}, q_{+ab}] + S_{BF \, eff., \, su(2)_{-}}[e^{I}_{\mu}, q_{-ab}] + \int V[e^{I}_{\mu}, q_{+ab}, q_{-ab}]$$

#### length metric and internal metrics from area metric

invert  $G(g, q_+, q_-) \longrightarrow g(G), q_+(G), q_-(G)$  (20 polynomial equations)

perturbative inversion in background expansion: [JB & Dittrich, 2022]

$$\begin{split} g_{\mu\nu} &= \delta_{\mu\nu} + h_{\mu\nu} \\ q_{\pm ab} &= \delta_{ab} + \chi_{\pm ab}, \quad \text{Tr}(\chi_{\pm}) = 0 \\ G_{\mu\nu\rho\sigma} &= 2\delta_{\mu[\rho}\delta_{\sigma]\nu} + a_{\mu\nu\rho\sigma}, \quad a_{\mu\nu\rho\sigma} = \mathbb{L}^{\lambda\tau}_{\mu\nu\rho\sigma}h_{\lambda\tau} + 2\mathbb{P}^{ab}_{+\mu\nu\rho\sigma}\chi_{+ab} + 2\mathbb{P}^{ab}_{-\mu\nu\rho\sigma}\chi_{-ab} \\ \text{invert } a(h, \chi_{+}, \chi_{-}) &\longrightarrow h(a), \chi_{+}(a), \chi_{-}(a) \\ S^{(2)}[a_{\mu\nu\rho\sigma}] \stackrel{?}{=} S^{(2)}_{P, mod. \, eff.}[h_{\mu\nu}, \chi_{+ab}, \chi_{-ab}] = S^{(2)}_{BF \, eff., \, su(2)_{+}}[h_{\mu\nu}, \chi_{+ab}] + S^{(2)}_{BF \, eff., \, su(2)_{-}}[h_{\mu\nu}, \chi_{-ab}] + \int V^{(2)}[\chi_{+ab}, \chi_{-ab}] \\ \end{split}$$



#### **quadratic potential introducing mass terms for the χ-fields:** [JB & Dittrich, 2022]

$$S_{P,mod.\,eff.}^{(2)}[h_{\mu\nu},\chi_{+\,ab},\chi_{-\,ab}] = \frac{\gamma_{+}}{2} \int \mathscr{L}_{EH}^{(2)}(h_{\mu\nu}+\chi_{+\,\mu\nu}) + \frac{\gamma_{-}}{2} \int \mathscr{L}_{EH}^{(2)}(h_{\mu\nu}+\chi_{-\,\mu\nu}) + \frac{\gamma_{+}}{8} \int m_{+}^{2}\chi_{+\,\mu\nu}\chi_{+}^{\mu\nu} + \frac{\gamma_{-}}{8} \int m_{-}^{2}\chi_{-\,\mu\nu}\chi_{-}^{\mu\nu}, \quad \gamma_{\pm} = \frac{\gamma_{+}}{2} \int \mathscr{L}_{EH}^{(2)}(h_{\mu\nu}+\chi_{+\,\mu\nu}) + \frac{\gamma_{-}}{2} \int \mathscr{L}_{EH}^{(2)}(h_{\mu\nu}+\chi_{-\,\mu\nu}) + \frac{\gamma_{-}}{8} \int m_{+}^{2}\chi_{+\,\mu\nu}\chi_{+}^{\mu\nu} + \frac{\gamma_{-}}{8} \int m_{-}^{2}\chi_{-\,\mu\nu}\chi_{-}^{\mu\nu}, \quad \gamma_{\pm} = \frac{\gamma_{+}}{2} \int \mathscr{L}_{EH}^{(2)}(h_{\mu\nu}+\chi_{+\,\mu\nu}) + \frac{\gamma_{-}}{2} \int \mathscr{L}_{EH}^{(2)}(h_{\mu\nu}+\chi_{-\,\mu\nu}) + \frac{\gamma_{-}}{8} \int m_{+}^{2}\chi_{+\,\mu\nu}\chi_{+}^{\mu\nu} + \frac{\gamma_{-}}{8} \int m_{-}^{2}\chi_{-\,\mu\nu}\chi_{-}^{\mu\nu}, \quad \gamma_{\pm} = \frac{\gamma_{+}}{2} \int \mathscr{L}_{EH}^{(2)}(h_{\mu\nu}+\chi_{+\,\mu\nu}) + \frac{\gamma_{-}}{2} \int \mathscr{L}_{EH}^{(2)}(h_{\mu\nu}+\chi_{+\,\mu\nu}) + \frac{\gamma_{-}}{8} \int m_{+}^{2}\chi_{+\,\mu\nu}\chi_{+}^{\mu\nu} + \frac{\gamma_{-}}{8} \int m_{-}^{2}\chi_{-\,\mu\nu}\chi_{-}^{\mu\nu}\chi_{-}^{\mu\nu}, \quad \gamma_{\pm} = \frac{\gamma_{+}}{2} \int \mathscr{L}_{EH}^{(2)}(h_{\mu\nu}+\chi_{+\,\mu\nu}) + \frac{\gamma_{-}}{2} \int \mathscr{L}_{EH}^{(2)}(h_{\mu\nu}+\chi_{+\,\mu\nu}) + \frac{\gamma_{-}}{8} \int m_{+}^{2}\chi_{+\,\mu\nu}\chi_{+}^{\mu\nu} + \frac{\gamma_{-}}{8} \int m_{-}^{2}\chi_{-\,\mu\nu}\chi_{-}^{\mu\nu}\chi_{-}^{\mu\nu},$$

$$\chi_{\pm \mu\nu} \equiv 4\mathbb{P}^{ab}_{\pm \mu\nu\rho\sigma} \frac{k^{\rho}k^{\sigma}}{\Delta} \chi_{\pm ab} \text{ isometric embedding } \{\chi_{ab} \mid \chi_{ab} = \chi_{ba}, \mathcal{I}_{ab} \in \mathcal{L}_{EH} \}$$

$$\mathscr{L}^{(2)}_{EH} \supset \text{ Pauli-Fierz operator } \mathscr{C}^{\mu\nu\rho\sigma} \equiv \frac{\Delta}{4} \left( 2P^{\mu\nu\rho\sigma} - 2^{0}P^{\mu\nu\rho\sigma} \right) \text{ (incl.)}$$

#### symmetries:

- linearized diffeomorphisms
- BF shift symmetry broken by mass terms propagating d.o.f. chiral:  $(h, \chi_{\pm}) \rightarrow (H_{\pm} \equiv h + \chi_{\pm}, \chi_{\pm})$  non-local field redefinition; integrate out  $\chi_{\pm} \longrightarrow S_{EH}^{(2)}[H_{\pm}]$  with "shifted graviton" non-chiral: existence & type of potential additional d.o.f. not settled

 $\operatorname{Tr}(\chi) = 0\} \to \{\chi_{\mu\nu} | \chi_{\nu\mu} = \chi_{\mu\nu}, \operatorname{Tr}(\chi) = 0\}, \ \chi_{ab}\chi^{ab} = \chi_{\mu\nu}\chi^{\mu\nu}, \ \partial^{\mu}\chi_{\mu\nu} = 0$ 

l. spin-0 & spin-2 projectors)





### **Effective length-metric action**

integrate out  $\chi_{\pm}$ : (in Fourier space,  $S^{(2)} \equiv \sum_{k} \mathscr{L}^{(2)}(k) = \sum_{k} \phi(-k)H(k)$ 

$$\mathscr{L}_{eff.}^{(2)}(h_{\mu\nu}) = \mathscr{L}_{EH}^{(2)}(h_{\mu\nu}) - \frac{1}{4}{}^{(1)}C_{\mu\nu\rho\sigma}\left(\frac{1}{m_{+}^{2} + \bigtriangleup} + \frac{1}{m_{-}^{2} + \bigtriangleup}\right)^{(1)}C^{\mu\nu\rho\sigma}$$

- GR to lowest order
- full corrections: quadratic in Weyl tensor, non-local, ghost-free (for  $m_{+} = m_{-}$ )

$$I(k)\phi(k), \quad \triangle = k_{\mu}k_{\nu}\delta^{\mu\nu}) \quad [JB \& Dittrich, 2022]$$

• lattice continuum limit of Area-Regge action reproduced at next order:  $\mathscr{L}_{eff.}^{(2)} = \mathscr{L}_{EH}^{(2)} + \text{Weyl}^2$  [Dittrich, 2021; Dittrich & Kogios, 2022]



### **Geometric interpretation of non-metric d.o.f.**

#### on-shell connection

Gauss law:  $d_{\omega}B = 0$ splits into two independent su(2) Gauss laws:  $d_{\omega_{\pm}}B^{i}_{\pm} = dB^{i}_{\pm} + \omega^{i}_{\pm j} \wedge B^{j}_{\pm} = 0$  (assuming non-degenerate tetrads)

#### two equivalent perspectives:

[Freidel, 2008; Speziale, 2010; Beke, Palmisano & Speziale, 2012]  $d_{\omega_{\pm}}B_{\pm}^{i} = 0 \quad \Leftrightarrow \quad d_{A_{\pm}}\Sigma_{\pm}(e) = 0, \quad d_{A_{\pm}}q_{\pm ab} = 0 \quad \text{where } A_{\pm} \equiv b_{\pm}^{-1}db_{\pm} + b_{\pm}^{-1}\omega_{\pm}b_{\pm} \quad \text{torsion-free & non-metric w.r.t } \delta_{ij}$ [JB & Dittrich, w.i.p.]  $d_{\omega_{\pm}}B_{\pm}^{i} = 0 \quad \Leftrightarrow \quad d_{\omega_{\pm}}\Sigma_{\pm}(e) \neq 0, \quad d_{\omega_{\pm}}\delta_{ij} = 0 \quad \text{i.e. } \omega_{\pm} \text{ torsion-full & metric w.r.t } \delta_{ij}$ 

#### Non-metric d.o.f. of area metric source torsion!



# **Summary & prospects**

- spin foam path integral defined on extended configuration space
- semi-classical configuration space parametrized by area metrics
- indications: continuum limit of Area-Regge action; geometry of a twisted simplex
- *G* area metric  $\leftrightarrow$  *g* length metric + scalars (—torsion)
- $S_{eff}[g] = S_{EH}[g] + Weyl^2$  non-local, ghost-free correction

### phenomenology:

- corrections to Newton potential & black hole entropy?
- hairy black holes à la Einstein-Weyl gravity & beyond?
- theoretical & observational bounds on scalar masses? constraints on  $\gamma$ ?

#### conceptual:

- full non-linear GR from area-metric dynamics? (via alternative parametrization of B-field?)
- canonical counting of d.o.f.
- different choices of potential

### fundamental:

- towards area metric geometry: area-metric compatible connection, curvature (continuum & discrete)
- "canonical" derivation of black hole entropy in language of area metrics?
- torsion scalar potential as dynamical or emergent cosmological constant? chiral vs non-chiral?
- possible relation to non-symmetric gravity & fundamental fields in string theory?

• continuum effective dynamics of spin foams reproduced by modified non-chiral Plebanski theories with area metric as configuration variable, obtained by imposing equal left-/right-handed length-metric geometries & replacing remaining half of simplicity constraints by potential





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