

els The future of spin foams



The good:

Areas are fundamental and have discrete spectrum.

(3D angles can be seen as auxiliary variables.)

Convergence with other quantum gravity approaches.

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Areas are fundamental and have discrete spectrum. (3D angles can be seen as auxiliary variables.)



We have not only metric degrees of freedom but additional (torsion) degrees of freedom.

We have an anomaly in the algebra of primary simplicity constraint and in the algebra of secondary simplicity constraints (gluing/ shape matching) parametrized by γ .

Convergence with other quantum gravity approaches.





The bad (?):

The anomaly is responsible for the "flatness problem".

[Asante, BD, Haggard 2020]

Applies to all models where areas are independent variables have approx. equidistant spectrum

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- b)

Numerical proof that this problem occurs in EPRL in "standard classical limit". [Gozzini 21]

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One can thus consider this issue in a model, which captures the above key ingredients of spin foams.

Effective spin foams models: Seconds on a laptop instead of weeks on HPC. Transparent encoding of the dynamics, in particular the constraints.

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In the discrete:

[Asante, BD, Haggard 20]

First proof that Regge equations of motion are reproduced for examples with inner edge. By computing the full non-perturbative partition function and expectation values of observables. Requires $\gamma \sim < 0.1$, but this allows for large deficit angles, consistent with Regge dynamics.

Concentrate on universal features

Consistent with semi-classical results by Han 2013, Han, Huang, Liu, Qu 2021: -bound on γ

- -complex saddle points with curvature in EPRL
- (as in effective spin foams)



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In the continuum:

Leveraging that non-perturbative results show good approximation by saddle point analysis near flat space: [BD 21] Perturbative analysis.

Study of linearized Area Regge action (+ $\frac{i}{v}$ constraints^2) on hypercubical lattice:

- leads to linearized Einstein Hilbert action (zeroth order in the lattice constant)
- a correction, which is of fourth order in the lattice constant, of six order in derivatives and quadratic in (derivative of) curvature

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- Different SF models differ in whether and how the non-metric degrees of freedom are suppressed.
- But are suppressed by dynamics in continuum limit anyway even without constraints.
- Emergence of universality in continuum limit.

Needed: "Effective" continuum action for spin foams (Area Regge action). Corrections to gravitational dynamics due to anomaly/ extended configuration space.

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• Non-metric degrees of freedom are getting very massive in the continuum limit: likely to extend to higher order perturbations





Concentrate on universal features for quantum gravity Lorentzian path integrals

Spin foams: one of the few approaches based on Lorentzian path integral. However, due to high numerical demands of EPRL/FK hardly explored issues important for Lorentzian path integrals. Using effective spin foams: Encountered issue of how to deal with integrations over infinite domains.

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- Key to understand non-perturbative continuum limit.
- **Euclidean configurations?** [Han, Liu]

• How to compute path integrals with highly oscillating amplitude and (practically) unbounded integration range? • What happens with the conformal factor problem of Euclidean quantum gravity (which killed almost all lattice approaches)?

• What configurations to sum over: Allow causally irregular configurations? Allow topology change in time? Appearance of

[BD, Gielen, Schander: Causally irregular configurations appear even in the simplest cosmological example.]



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- Key to understand non-perturbative continuum limit.
- **Euclidean configurations?** [Han, Liu]

- [Asante, BD, Jia, Padua-Arguelles: to appear]
- Picard Lefshetz methods
- suppressed.

There is a lot to explore for Lorentzian path integrals. Need effective numerical models and methods.

• How to compute path integrals with highly oscillating amplitude and (practically) unbounded integration range? • What happens with the conformal factor problem of Euclidean quantum gravity (which killed almost all lattice approaches)?

• What configurations to sum over: Allow causally irregular configurations? Allow topology change in time? Appearance of

[BD, Gielen, Schander: Causally irregular configurations appear even in the simplest cosmological example.]

[Feldbrugge, Lehners Turok; Han, Huang, Liu, Qu, Wan; Asante, BD, Jia, Padua-Arguelles: to appear]

• In simplest examples: make integrals (quickly) convergent. Conformal factor rotated to suppressing Euclidean branch. • Can allow for causally irregular configurations: With Picard-Lefshetz methods - indications that these will be always

> Convergence with Causal Dynamical Triangulations: causally irregular configurations suppressed Convergence with results from quantum cosmology with PL-integration

> > [Asante, BD, Padua-Arguelles 2021]







Universal Summary

- Concentrate on universal features of spin foams.
- E.g. the Lorentzian path integral.

Need effective numerical models and methods.



Areas as fundamental degrees of freedom.

[Asante, Bahr, BD, Steinh • Amplitudes of any model will change under coarse graining flow: universal features will survive.

• Concentrate on universal features of quantum gravity. Consider issues which are also of interest to other approaches.

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Thank you!



Effective Spin Foam Models

Effective Spin Foam models:

- 3D angles are already integrated out.
- much more transparent encoding of the dynamics, in particular with regard to gluing constraints
- form can be motivated from higher gauge theory (as opposed to gauge theory)

$$Z = \sum_{\text{discr.areas}} \exp\left(\frac{i}{\ell_P^2} S_{\text{Regge}}(\{A_t\})\right) G(\{A_t\}),$$

Oscillating factor with Area Regge action
(motivated by higher gauge theory)
Gaussian factor peaked on constraints.
Deviation proportional to $\ell_P \sqrt{\gamma}$ Area
(Determined from anomaly/ non-commutativity.)



[Asante, BD, Haggard PRL 2020]

• much much more amenable to numerical investigations: seconds on laptop compared to weeks on HPC

The number of angle variables is $5 \times \#(4\text{-simplices})$. Integrating them out beforehand saves a lot of computational resources.

Naive $\ell_P \to 0$ limit: Oscillations win over Gaussian. Constraints are not 'visible'. Leading to flatness problem for spin foams.

To avoid washing out of Gaussians we need (via naive estimation):

$$\sqrt{\frac{\sqrt{\text{Area}}}{\ell_P}} \text{curv}_t \le \mathcal{O}(1)$$

[Han, Asante-BD-Haggard]

A proper semi-classical regime does also require a small Barbero-Immirzi parameter. (It indeed is an anomaly parameter.)











Explicit computation of partition functions and expectation values for small triangulations

Triangulations with inner edge:

Computed via an exact evaluation of the path integral with observable insertions expectation values of various geometric observables.

We showed that we do have a regime of scale, γ and curvature, for which the correct equations of motions are implemented.

First explicit proof that spin foams can implement the correct equations of motions. (With finite, non-vanishing curvature angles, finite but small γ .)



Discrete dynamics: the Barbero-Immirzi parameter is not a free parameter.

Explicit tests

[Asante, BD, Haggarrd CQG 20]

Found that semi-classical regime can be even larger than suggested by naive bound, in particular if curvature (per triangle) is small.

Instabilities in expectation values: non-perturbative effects, resulting from interplay of discrete spectra and constraint implementation.

For all tested curvatures: $\gamma \sim 0.1$ defines semi-calssical regime.





Scaling of the Hessian

• λ and Λ both give the lattice constant, but originate from different sources. Λ counts derivatives.

$$H_{(ab)(ab)} \sim \frac{1}{\lambda^2} \times \frac{(a_1)}{(a_2)} (a_2)}{(a_1)} (a_2) (b_1) (b_2)}{(a_1)} \times \frac{(a_1)}{(a_2)} \Lambda^2 0 \Lambda^3 0}{(a_2)} \frac{\Lambda^0}{\Lambda^0} \Lambda^1 \Lambda^1}{(b_1)} \frac{\Lambda^0}{\Lambda^3} (A^1) \Lambda^0}{(b_2)} \frac{(a_2)}{(a_2)} 0 \Lambda^1 0 \Lambda^0}{(a_2)} \frac{\Lambda^0}{\Lambda^0}$$

• With this scaling we already know that:

$$\Gamma_{(a)(a)} \sim \frac{1}{\lambda^2} \times \frac{(a_1)(a_2)}{(a_1)\Lambda^6} \frac{\Lambda^4}{\Lambda^4}}{(a_2)\Lambda^4}$$

• Correction from integrating out the zeta-variables.

$$\Gamma_{(a)(a)} = -H_{(a)(b)} \cdot H_{(b)(b)}^{-1} \cdot H_{(b)(b)}$$

• $H_{(b)(b)}$ is invertible (perturbatively without inverse derivatives)! (This is where we might have needed the G-matrix: but it is not necessary here!)

• That is the correction for metric sector scales with a_L^4 . • Integrating out the spurious (a2) variables leads to even higher order corrections.

