Causal Structure in Spin-foams

Eugenio Bianchi & <u>Pierre Martin-Dussaud</u> arXiv: 2109.00986

ILQGS, 22 Feb. 2022





THE QUANTUM INFORMATION STRUCTURE OF SPACETIME



Causal Structure in Spin-foams

Is causality a fundamental or an emergent property of space-time?



Causal Structure in Spin-foams

Is causality a fundamental or an emergent property of space-time?

<u>Plan:</u>

kinematics of general relativity

+ discrete

+surface d.o.f.

+ dynamics

+quantum: path-integral spinfoams (crash-course)

causal spinfoams



manifold + lorentzian metric g



manifold + lorentzian metric g

signature
$$(\eta, -\eta, -\eta, -\eta)$$
 $\eta = \pm 1$



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causal structure = light-cone structure



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causal structure = light-cone structure

= equivalence classes on T_xM

$$\operatorname{sign} g(u, u) = \begin{cases} \eta & \operatorname{Time-like} \\ 0 & \operatorname{Null} \\ -\eta & \operatorname{Space-like} \end{cases}$$



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No time arrow!



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No time arrow!

Malament's theorem (1977): Metric = Causal structure + Conformal factor $\Omega(x)$





2-dimensional differentiable manifold





2-dimensional differentiable manifold

complex of triangles





2-dimensional differentiable manifold

complex of triangles



Triangle Segment Point







2-dimensional differentiable manifold

complex of triangles

dual skeleton



Triangle Segment Point Vertex Link Face



4-dimensional differentiable manifold with lorentzian metric





4-dimensional differentiable manifold with lorentzian metric

complex of 4-simplices (with space-like tetrahedra)





4-dimensional differentiable manifold with lorentzian metric

complex of 4-simplices (with space-like tetrahedra)



4-simplex Tetrahedron Triangle Segment Point



4-dimensional differentiable manifold with lorentzian metric

complex of 4-simplices (with space-like tetrahedra)



4-simplex Tetrahedron Triangle Segment Point

dual skeleton with local time arrows Vertex Link Face ••• ...



4-dimensional differentiable manifold with lorentzian metric



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+surface d.o.f.







$$arepsilon_v(e) = \left\{ egin{array}{c} -1 & {
m incoming (past)} \\ 1 & {
m outgoing (future)} \end{array}
ight.$$



+surface d.o.f.

dual 1-skeleton



Causal constraint

 $\varepsilon_v(e) = \begin{cases} -1 & \text{incoming (past)} \\ 1 & \text{outgoing (future)} \end{cases}$



+surface d.o.f.



$$\varepsilon_v(e) = \left\{ \begin{array}{cc} -1 & \text{incoming (past)} \\ 1 & \text{outgoing (future)} \end{array} \right.$$

$$\varepsilon_v(f) = \left\{ \begin{array}{cc} \eta & \text{co-chronal (time)} \\ -\eta & \text{anti-chronal (space)} \end{array} \right.$$











 $\varepsilon_v(f) = \begin{cases} \eta & \text{co-chronal (time)} \\ -\eta & \text{anti-chronal (space)} \end{cases}$

+surface d.o.f.

The causal structure of a 2-complex can be encoded by assigning an orientation (a sign) either to the edges, or to the wedges, satisfying some "causal constraints".

 $\varepsilon_v(f) = \eta \varepsilon_v(e_1) \varepsilon_v(e_2)$

1:2

No time-arrow





 $\varepsilon_v(e) = \begin{cases} -1 & \text{incoming (past)} \\ 1 & \text{outgoing (future)} \end{cases}$

+dynamics



Palatini action

$$S(g_{\mu\nu},\Gamma) = \int d^4x \sqrt{-g} g^{\mu\nu} R_{\mu\nu}(\Gamma)$$

[Barrett 94']

 1^{st} order Regge action

+dynamics

 $S[l_s, \theta_{t\sigma}, \mu_{\sigma}] = \sum \sum A_t(l_s)\theta_{t\sigma} + \sum \mu_{\sigma} \det \gamma_{\sigma}$ σ $t|t\in\sigma$ σ





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sum over 4-simplices

sum over wedges



Palatini action

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\mathcal{M}

Palatini action

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+quantum: path integral

$$\mathcal{A}(l_{\Sigma},\theta_{\Sigma}) = \int [dl_s] [d\theta_{t\sigma}] \prod_{\sigma} \delta(\det \gamma_{\sigma}) e^{\frac{i}{\hbar} \sum_{t \in \sigma} A_t \theta_{t\sigma}}$$



+quantum: path integral $\varepsilon = \pm 1$ a > 0

$$\mathcal{E} = \pm 1 \quad \rho > 0$$
$$\mathcal{E} \rho$$
$$\mathcal{A}(l_{\Sigma}, \theta_{\Sigma}) = \int [dl_s] [d\theta_{t\sigma}] \prod_{\sigma} \delta(\det \gamma_{\sigma}) e^{\frac{i}{\hbar} \sum_{t \in \sigma} A_t \theta_{t\sigma}}$$



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1. This sum is made over all possible configurations of ε assigned to wedges, including both causal and non-causal (spurious) configurations.

$$\prod_{f \in \text{cycle}} \varepsilon_v(f) = \eta^{\text{\#cycle}}$$



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 The presence of non-causal histories contributing to the amplitude depends on a certain choice of variables.
 Example: compare with standard Regge calculus.



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4. The choice of including the spurious terms or not depends on what we want to compute: a projector on the physical Hilbert space or a causal propagator? Feynman propagator: $(\Box + m^2) W = \delta$ Hadamard function: $(\Box + m^2) W = 0$ [Teitelboim

[Teitelboim 82'] [Livine, Oriti 02']





Spin-network:

(EPRL)



Spin-network:

(EPRL)

Transition amplitude

 $W(s_i, s_f) = \sum_{\mathcal{C}} W_{\mathcal{C}}(s_i, s_f)$ sum over 2-complexes



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Spin-foam amplitude $W_{\mathcal{C}}(s_i, s_f) = \int_{SU(2)} [dh_w] \prod_{f} \delta(\prod_{w \in f} h_w) \prod_{v} A_v(h_w)$ wedge variable \in SU(2)

Spin-network:

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Spin-network:

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Spin-foam amplitude $W_{\mathcal{C}}(s_i, s_f) = \int_{SU(2)} [dh_w] \prod_{f} \delta(\prod_{w \in f} h_w) \prod_{v} A_v(h_w)$ wedge variable \in SU(2) Vertex amplitude $A_{v}(h_{w}) = \int_{SL_{2}(\mathbb{C})} [dg_{e}] \prod_{w \in v} K(h_{w}, g_{s_{w}}g_{t_{w}}^{-1})$ edge variable $\in SL_{2}(\mathbb{C})$ wedges

Spin-network:

quantum geometry of space

Wedge amplitude

$$K(h,g) = \sum_{j} \frac{(2j+1)^4}{\pi^3} \int [d\zeta] [dz'] [dz''] \langle z' | h^{\dagger} | z'' \rangle^{2j} \ \mathcal{B} e^{iS_{\gamma}}.$$





$$\begin{array}{l} \begin{array}{c} \mathsf{Causal spinfoams} \\ \text{Wedge amplitude} \\ K(h,g) = \sum_{j} \frac{(2j+1)^4}{\pi^3} \int_{\text{integrations over CP}^1} |d\zeta| |dz''| \langle z'|h^{\dagger}|z''\rangle^{2j} & \mathcal{B} e^{iS\gamma}. \end{array} \\ \begin{array}{c} \text{sum over spins} \\ \text{Sum over spins} \end{array} \\ \begin{array}{c} \mathcal{S}_{\gamma} = \gamma j \log \frac{\|g^T \zeta\|^2}{\|\zeta\|^2} \\ \text{semi-classical limit} \\ \theta_w \end{array} \\ \begin{array}{c} \mathcal{S}_{\gamma} = \gamma j \log \frac{\|g^T \zeta\|^2}{\|\zeta\|^2} \\ \text{semi-classical limit} \\ \theta_w \end{array} \\ \begin{array}{c} \mathcal{S}_{\gamma} = \gamma j \log \frac{\|g^T \zeta\|^2}{\|\zeta\|^2} \\ \text{semi-classical limit} \\ \theta_w \end{array} \\ \begin{array}{c} \mathcal{S}_{\gamma} = \gamma j \log \frac{\|g^T \zeta\|^2}{\|\zeta\|^2} \\ \text{semi-classical limit} \\ \theta_w \end{array} \\ \begin{array}{c} \mathcal{S}_{\gamma} = \gamma j \log \frac{\|g^T \zeta\|^2}{\|\zeta\|^2} \\ \text{semi-classical limit} \\ \theta_w \end{array} \\ \begin{array}{c} \mathcal{S}_{\gamma} = \gamma j \log \frac{\|g^T \zeta\|^2}{\|\zeta\|^2} \\ \text{semi-classical limit} \\ \theta_w \end{array} \\ \end{array} \\ \begin{array}{c} \mathcal{S}_{\gamma} = \gamma j \log \frac{\|g^T \zeta\|^2}{\|\zeta\|^2} \\ \frac{\mathcal{S}_{\gamma}}{\|g\|^2} \\ \frac{\mathcal{S}_{\gamma}}$$

$$\begin{aligned} & \underset{\text{Wedge amplitude}}{\text{Sum over spins}} & \underset{j}{\text{(EPRL)}} & \underset{\text{integrations over CP}^{\text{(eprl)}}}{\text{(EPRL)}} & \underset{\text{wedge action}^{\text{``wedge action}^{\text{``}}}}{\int G(z) [dz'] [dz''] (dz''] (dz'') (dz'') (dz'') (dz'')} & \underset{\text{semi-classical limit}}{\int G(z)} & \underset{\text{semi-classical limit}}{\int G(z)} & \underset{\text{wedge action}^{\text{``wedge action}^{\text{``}}}}{\int G(z) [dz''] (dz'') (dz''') (dz'') (dz'') (dz'') (dz'') (dz'') (dz''') (dz'') (dz'') (dz''$$

$$\begin{aligned} & \underset{\text{(EPRL)}}{\text{Wedge amplitude}} & \underset{\text{(EPRL)}}{\text{(EPRL)}} & \underset{\text{(wedge action"}}{\text{(sum over spins)}} & S_{\gamma} = \gamma j \log \frac{\|g^{T} \zeta\|^{2}}{\|\zeta\|^{2}} \\ & K(h,g) = \sum_{j} \frac{(2j+1)^{4}}{\pi^{3}} \int [d\zeta][dz'][dz''] \langle z'|h^{\dagger}|z''\rangle^{2j} & \mathcal{B} e^{iS_{\gamma}}. & \text{semi-classical limit} \\ & \theta_{w} \end{aligned} \\ & K(h,g) = \sum_{\varepsilon \in \{-1,1\}} K^{\varepsilon}(h,g) \\ & K^{\varepsilon}(h,g) = \sum_{j} \frac{(2j+1)^{4}}{\pi^{3}} \int [d\zeta][dz'][dz''] \langle z'|h^{\dagger}|z''\rangle^{2j} & \mathcal{B} \Theta(\varepsilon S_{\gamma}) e^{iS_{\gamma}}. \end{aligned} \\ & \text{Vertex amplitude} \\ & A_{v} = \sum_{\substack{[\varepsilon_{w}] \\ \text{causal } \eta = 1}} A_{v}^{\varepsilon} + \sum_{\substack{[\varepsilon_{w}] \\ \text{causal } \eta = -1}} A_{v}^{\varepsilon} + \sum_{\substack{[\varepsilon_{w}] \\ \text{spurious}}} A_{v}^{\varepsilon} \end{aligned}$$

 $\stackrel{\rm signature}{(\eta,-\eta,-\eta,-\eta)}$



Causal spinfoams

(EPRL)

Spin-foam amplitude



Causal spinfoams

Spin-foam amplitude



Causal spinfoam: propagator of spin-networks (time-evolution)

Full EPRL spinfoam: projector on the physical states (Hamiltonian constraint)

Causal spinfoams

[Christodoulou, Längvik, Riello, Röken, Rovelli, 13']

Spin-foam amplitude $W_{\mathcal{C}} = \sum_{\substack{[\varepsilon_w]\\\text{causal }\eta=1}} W_{\mathcal{C}}^{\varepsilon} + \sum_{\substack{[\varepsilon_w]\\\text{causal }\eta=-1}} W_{\mathcal{C}}^{\varepsilon} + \sum_{\substack{[\varepsilon_w]\\\text{spurious}}} W_{\mathcal{C}}^{\varepsilon}$

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[Livine, Oriti 03']

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[Immirzi, Gupta, Smolin, Markopoulou...]



Causal spinfoam: propagator of spin-networks (time-evolution)

Full EPRL spinfoam: projector on the physical states (Hamiltonian constraint)

Livine, Oriti 03'] For the Barret-Crane model:

$$K^{p}(x_{1}, x_{2}) = \frac{2\sin(\beta(x_{1}, x_{2}) p/2)}{p\sinh\beta(x_{1}, x_{2})} = \frac{1}{p\sinh\beta(x_{1}, x_{2})} \sum_{\varepsilon = \pm 1} \varepsilon e^{i\varepsilon\beta(x_{1}, x_{2}) p/2}$$

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[Engle 13'] "Proper vertex amplitude"

$$A_v^{(+)}(\{k_{ab},\psi_{ab}\}) \coloneqq (-1)^{\Xi} \int_{\mathrm{SL}(2,\mathbb{C})^5} \delta(X_4) \prod_a \mathrm{d}X_a \prod_{a < b} \alpha(X_a \mathcal{I}\psi_{ab}, X_b \mathcal{I} \prod_{ba} (\{X_{ab}\}) \psi_{ba})$$

[Immirzi, Gupta, Smolin, Markopoulou...]



Conclusion

- Causality exhausts almost all of the information content of the gravitational field (Malament's theorem).
- Causality can be represented by arrows on the edges of the dual skeleton.
- Causality can also be encoded on the wedges (more variables) provided a new "causal constraint" is satisfied.
- This causal constraint can be derived from the equations of motion.
- In the path-integral, one can restrict or not to causal histories only, depending on what one wants to compute, a projector or a propagator.
- Same thing with spinfoams.

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