LQC of black holes (recent progress)



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Beatriz Elizaga Navascués, Louisiana State University

- Black holes: Promising observational windows to quantum gravity phenomena, in the dawn of GW astronomy.
- Long history of works for the quantization of black holes in LQG & LQC. Satisfactory yet simple description?

[Ashtekar, Bojowald, Modesto, Cartin, Khanna, Boehmer, Vandersloot, Chiou, Campiglia, Gambini, Pullin, Sabharwal, Brannlund, Kloster, De Benedictis, Olmedo, Dadhich, Joe, Singh, Haggard, Rovelli, Corichi, Saini, Cortez, Cuervo, Morales-Técotl, Ruelas, Yonika, Bianchi, Christodoulo, D'Ambrosio, Alesci, Bahrami, Pranzetti, Kelly, Santacruz, Wilson-Ewing, Zhang, Ma, Song,
Bodendorfer, Mele, Münch, Mena Marugán, García-Quismondo, Perez, Speziale, Viollet, Han, Liu, Alonso-Bardaji, Brizuela, Vera, ...]

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- Long history of works for the quantization of black holes in LQG & LQC. Satisfactory yet simple description?
 - ★ Spherical symmetry as starting point.
 - ★ Use LQC techniques to describe the interior.

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- Highlight: Extension of Kruskal spacetime by AOS. [Ashtekar, Olmedo, Singh 2019]

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★ Quantization? (General properties, eff. regimes, etc.)



BH interior & its "effective" LQC models

Interior region in GR

• Metric: Cosmology of Kantowski-Sachs type.

$$ds^{2} = -N(t)^{2}dt^{2} + \frac{p_{b}^{2}(t)}{L_{o}^{2} |p_{c}(t)|} dx^{2} + |p_{c}(t)| d\Omega_{S}^{2}, \quad x \in (0, L_{o})$$

• Symmetry-reduced Ashtekar-Barbero variables:

$$E_i^{\alpha} \to (p_b, p_c) \qquad A_{\alpha}^{i} \to (b, c)$$
$$\{b, p_b\} = \gamma, \qquad \{c, p_c\} = 2\gamma, \qquad \gamma \to \text{Immirzi}$$

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$$H_{KS}[N] = NL_o \frac{b}{\gamma \sqrt{|p_c|}} \left(O_b^{KS} - O_c^{KS} \right) \longrightarrow O_b^{KS} = O_c^{KS} = m$$

 \propto ADM mass

Regularization of Hamiltonian

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- Regularized Hamiltonian [ignoring $\mathcal{O}(\delta_b)$]:

 $H_{KS}^{\text{reg}}[N_T] = L_o O_b^{KS} \left[b \to \sin(\delta_b b) / \delta_b \right] - L_o O_c^{KS} \left[c \to \sin(\delta_c c) / \delta_c \right],$

 $N_T = \frac{\gamma \delta_b \sqrt{|p_c|}}{\sin(\delta_b b)}, \quad |\delta_b|, |\delta_c| \text{ coord. lengths of min. area plaquettes}$

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"Effective" models

$$H_{KS}^{\text{reg}}[N] = NL_o \frac{\sin(\delta_b b)}{\gamma \delta_b \sqrt{|p_c|}} \left(O_b - O_c \right), \quad \text{effLQC} : O_b = O_c = m$$

- Several studies of the evolution generated by $H_{KS}^{reg}[N]$.
- Freedom in the choice of parameters δ_b, δ_c .

[Ashtekar, Bojowald, Modesto, Campiglia, Gambini, Pullin, Boehmer, Vandersloot, Chiou, Joe, Singh, Corichi, Olmedo, Saini, Bodendorfer, Mele, Münch, Mena Marugán, García-Quismondo, Han, Liu, ...]

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[Ashtekar, Olmedo, Singh]

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- 2019: "Quantum Transfiguration of Kruskal BHs" (AOS),

$$\delta_b \sim \left(\frac{\sqrt{\Delta}}{\sqrt{2\pi\gamma^2 m}}\right)^{1/3}, \quad L_o \delta_c \sim \frac{1}{2} \left(\frac{\gamma \Delta^2}{4\pi^2 m}\right)^{1/3}, \quad \text{on solution}$$



AOS type of models: Hamiltonian formulation

Caveats on derivation of e.o.m

• The AOS "effective" dynamics is generated by:

$$H_{\text{AOS}}^{\text{eff}} = L_o \left(O_b - O_c \right) \left(= H_{KS}^{\text{reg}}[N_T] \right)$$

treating the parameters δ_b, δ_c as pure constants.

- On each solution, constraint implies $O_b = O_c = m$.
- On each sol., declare $\delta_b = \delta_b(m)$, $\delta_c = \delta_c(m)$ for large BHs.

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- On each sol., declare $\delta_b = \delta_b(m)$, $\delta_c = \delta_c(m)$ for large BHs.
- However, *m* is a function on phase space (const. motion)!
 [Bodendorfer, Mele, Münch, Mena Marugán, García-Quismondo]

Extended phase space

• The AOS model can be derived from a Hamiltonian

$\underline{N}H_{AOS}^{\text{eff}} + \lambda_b \Psi_b + \lambda_c \Psi_c,$

on a extended phase space that includes δ_b , δ_c (& momenta).

• Ψ_b, Ψ_c : Relation between δ_b, δ_c and BH mass on solutions:

$$\Psi_b = K_b \left(O_b, O_c \right) - \delta_b, \quad \Psi_c = K_c \left(O_b, O_c \right) - \delta_c$$

$$K_b(m,m) \sim \left(\frac{\sqrt{\Delta}}{\sqrt{2\pi\gamma^2 m}}\right)^{1/3}, \quad K_c(m,m) \sim \frac{1}{2L_o} \left(\frac{\gamma \Delta^2}{4\pi^2 m}\right)^{1/3}$$

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• AOS dynamics follows for $\lambda_b = 0 = \lambda_c$.

Gauge reduction

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- <u>Idea</u>: fix the gauge so that $\lambda_b = \lambda_c = 0$ to remove δ_b, δ_c .
- <u>Result</u>: The reduced phase space is <u>**not**</u> symplectomorphic to that of Kantowski-Sachs cosmologies (≠ GR algebra).

[BEN, Mena Marugán, García-Quismondo]

Road to quantization

• The AOS model can be derived from a Hamiltonian

$\underline{N}H_{AOS}^{eff} + \lambda_b \Psi_b + \lambda_c \Psi_c,$

on a extended phase space that includes δ_b , δ_c (& momenta).

• LQC under control for triad/connection algebra in GR.

• Quantization strategy: Use LQC tools to find a quantum representation of constraints in the extended framework.



Quantization: Kinematics

Holonomy-flux algebra

• Matrix elements of holonomies of A^i_{α} along θ and x:

$$\mathcal{N}_{\mu_b} = e^{ib\mu_b/2}, \qquad \mathcal{N}_{\mu_c} = e^{ic\mu_c/2}, \qquad \mu_b, \mu_c \in \mathbb{R}.$$

- Fluxes of E_i^{α} across surfaces determined by p_b and p_c .
- Holonomy-flux algebra:

$$\{\mathcal{N}_{\mu_b}, p_b\} = \frac{i}{2} \mu_b \gamma \mathcal{N}_{\mu_b},$$

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Two copies of homogeneous & isotropic LQC

Kinematic Hilbert space

• \mathscr{H}_{LOC}^{kin} : space of square summable functions on \mathbb{R}^2 .

Basis $|\mu_b, \mu_c\rangle$ $\langle \mu_b, \mu_c | \mu'_b, \mu'_c, \rangle = \delta_{\mu_b, \mu'_b} \delta_{\mu_c, \mu'_c}$

• $\hat{\mathcal{N}}_{\mu_{b,c}} \rightarrow$ translation, $\hat{p}_{b,c} \rightarrow$ multiplication operators.

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- $\hat{\mathcal{N}}_{\mu_{b,c}} \rightarrow$ translation, $\hat{p}_{b,c} \rightarrow$ multiplication operators.
- Schrödinger representation for δ_b , δ_c and momenta.
- Total (kinematic) Hilbert space:

$$\mathcal{H}_T^{\mathrm{kin}} = \mathcal{H}_{\mathrm{LQC}}^{\mathrm{kin}} \otimes L^2(\mathbb{R}, d\delta_b) \otimes L^2(\mathbb{R}, d\delta_c)$$

$$H_{\rm AOS}^{\rm eff} = L_o \left(O_b - O_c \right)$$

• The "angular" part reads

$$O_{c} = \frac{\sin(\delta_{c}c)}{\gamma L_{o}\delta_{c}} p_{c}, \qquad \sin(\delta_{c}c) = \frac{1}{2i} \left(\mathcal{N}_{2\delta_{c}} - \mathcal{N}_{-2\delta_{c}} \right)$$

• Quantization: $\sin(\delta_c c) p_c / \delta_c \to \hat{\Omega}_c$ on (a subset of) $\mathscr{H}_T^{\text{kin}}$.

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- For each fixed pair δ_b , δ_c change the LQC basis:

$$|\mu_b,\mu_c\rangle \to |\tilde{\mu}_b = \mu_b \delta_b^{-1}, \tilde{\mu}_c = \mu_c \delta_c^{-1}\rangle$$

 $\hat{\Omega}_{c}^{2}$ same as in FLRW cosmology

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• Operator $\hat{O}_c \propto \hat{\Omega}_c$ essentially self-adjoint on

 $\operatorname{Cyl}_{\varepsilon_{c}}^{\pm} = \operatorname{span}\{ |\tilde{\mu}_{b}, \tilde{\mu}_{c}\rangle : \tilde{\mu}_{c} = \pm (\varepsilon_{c} + 2n), n \in \mathbb{N} \}, \quad \varepsilon_{c} \in (0, 2]$

• Abs. continuous spectrum equal to \mathbb{R} , independent of δ_c .

[Lewandowski, Kaminski, Mena Marugán, Martín-Benito, Olmedo, Pawlowski]

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$$\operatorname{Cyl}_{\varepsilon_c}^{\pm} \leftrightarrow \operatorname{span}\{ | \tilde{\mu}_b \in \mathbb{R} \rangle \otimes L^2(\mathbb{R}, dm) \}$$

Hamiltonian constraint $H_{AOS}^{eff} = L_o \left(O_b - O_c \right)$

• The "radial" part reads

$$O_b = -\frac{1}{2\gamma L_o} \left[\Omega_b + \gamma^2 \delta_b^2 \tilde{p}_b^2 \Omega_b^{-1} \right], \quad \Omega_b = \frac{\sin(\delta_b b)}{\delta_b} p_b, \quad \tilde{p}_b = \frac{p_b}{\delta_b}$$

• Quantization: "Minimal" requirements on \hat{O}_b .

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- Quantization: "Minimal" requirements on \hat{O}_b :
 - ★ Built from $\hat{\Omega}_b, \hat{\tilde{p}}_b$: Only depends on δ_b^2 for each pair δ_b, δ_c .
 - * Preserves super-selection sectors of $\hat{\Omega}_b$ (similar to $\hat{\Omega}_c$).
 - \star Ess. self-adjoint, no singular sp., point sp. = discrete.



Physical states

(in collab. with A. García-Quismondo, G.A. Mena Marugán & A. Mínguez)

Constraint equations

- Total Hamiltonian: $\underline{N}H_{AOS}^{eff} + \lambda_b \Psi_b + \lambda_c \Psi_c$.
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 $\hat{O}_b \psi_p(\tilde{\mu}_b, m, \delta_b, \delta_c) = m \psi_p(\tilde{\mu}_b, m, \delta_b, \delta_c), \quad m \in \mathbb{R}$

 $\psi_p = 0 \text{ for } m \notin \mathrm{Sp}_{\hat{O}_b}[\delta_b]$

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Looking for physical states in the algebraic dual of: $C_0^{\infty}(\operatorname{Sp}_{\hat{O}_b}^c) \cup \operatorname{span}\{ | \omega_n \in \operatorname{Sp}_{\hat{O}_b}^d \rangle \}$

Physical Hilbert space

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- Additional constraints: $\delta_b = K_b(O_b, O_c), \ \delta_c = K_c(O_b, O_c).$
- Quantum imposition of all constraints on $(\psi_p | :$

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• Inner product on \mathscr{H}_p from spectral measure of \hat{O}_b .

Possibilities: Discrete mass

$$O_b = -\frac{1}{2\gamma L_o} \left[\Omega_b + \gamma^2 \delta_b^2 \Omega_b^{-1} \tilde{p}_b^2 \right], \quad \Omega_b = \frac{\sin(\delta_b b)}{\delta_b} p_b, \quad \tilde{p}_b = \frac{p_b}{\delta_b}$$

• Imagine that it is possible to define \hat{O}_b such that

 $\hat{O}_b^2 \propto \hat{\Omega}_b^2 + 2\gamma^2 \delta_b^2 \hat{\tilde{p}}_b^2 + \gamma^4 \delta_b^4 \hat{P}, \qquad \hat{P} \text{ positive}$

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 $\langle \hat{O}_b^2 \rangle \ge \propto 2\gamma^2 \delta_b^2 \langle \hat{\tilde{p}}_b^2 \rangle \longrightarrow \operatorname{Sp}_{\hat{O}_b^2}[\delta_b \neq 0] \text{ discrete}$

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• Then $\text{Sp}_{\hat{O}_b}[\delta_b \neq 0]$ discrete: Gaps in the physical BH mass.

[Zhang, Ma, Song, Zhang, 2020-2022]

- Consistent with recent works quantizing $H_{AOS}^{eff} \sin(\delta_b b) p_b / \delta_b$.
- Satisfactory $\delta_b, \delta_c \to 0$ limit? Spectrum becomes continuous.

Possibilities: Cont. mass $O_b = -\frac{1}{2\gamma L_o} \left[\Omega_b + \gamma^2 \delta_b^2 \Omega_b^{-1} \tilde{p}_b^2 \right], \quad \Omega_b = \frac{\sin(\delta_b b)}{\delta_b} p_b, \quad \tilde{p}_b = \frac{p_b}{\delta_b}$

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- Imagine (e.g.) a "direct" quantization of O_b .
- Quantum constraint $(\psi_p | \hat{O}_b = (\psi_p | \hat{O}_c \text{ may be understood as})$

$$(\tilde{\psi}_p | (\hat{\Omega}_b^2 + \gamma^2 \delta_b^2 \hat{\tilde{p}}_b^2) = - (\tilde{\psi}_p | 2\gamma L_o m \hat{\Omega}_b \quad \text{with } (\tilde{\psi}_p | = (\psi_p | \hat{\Omega}_b^{-1}))$$

• Experience in LQC indicates that this eq. is solvable for continuous values of the BH mass. Good $\delta_b \rightarrow 0$ limit.

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- The standard habitat for physical states may be too small.
- Imagine (e.g.) a "direct" quantization of O_b .
- Quantum constraint $(\psi_p | \hat{O}_b = (\psi_p | \hat{O}_c \text{ may be understood as})$

$$(\tilde{\psi}_p | (\hat{\Omega}_b^2 + \gamma^2 \delta_b^2 \hat{\tilde{p}}_b^2) = - (\tilde{\psi}_p | 2\gamma L_o m \hat{\Omega}_b \quad \text{with } (\tilde{\psi}_p | = (\psi_p | \hat{\Omega}_b^{-1}))$$

Now we search on the algebraic

dual of $D(\hat{O}_{h})$ ("much" bigger!!)



Conclusions & outlook

- Aim: LQC of a BH interior leading to AOS-type models.
- Extended phase space: motivation from dynamics of effective model & convenient for quantization.
- Operators in Hamiltonian from FLRW LQC.
- <u>Important</u>: Choice of dual space to construct physical Hilbert space → Discrete/continuous BH mass.
- WdW limit? Relational dynamics? Effective AOS?