### Quantum deformation of 4d spin foam models

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### Introduction

- Spin foam models : discretised functional integrals for TFTs of BF type and quantum gravity
- Models for theories with zero 'cosmological constant' are based on the representation theory of Lie groups
- Problem : generally these models diverge
- Natural regularisation : consider models based on the representation theory of quantum groups
- Our work : application of this procedure to the EPRL model

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#### Outline

Three-dimensional models Topological models in four dimensions Quantum gravity models in four dimensions

#### Three-dimensional models

- Gravity in 2+1 dimensions
- The Ponzano-Regge model
- The Turaev-Viro model

#### 2 Topological models in four dimensions

- BF theory
- The Ooguri model
- The Crane-Yetter model

#### Quantum gravity models in four dimensions

- Gravity as a constrained topological theory
- The EPRL/FK models
- Quantum deformation of the EPRL model

Gravity in 2 + 1 dimensions The Ponzano-Regge model The Turaev-Viro model

# Gravity in 2 + 1 dimensions

- First order gravity on a 3-dimensional manifold *M*:
  - dual co-frame  $e: TM \to \mathbb{R}^3$
  - SO( $\eta$ ) connection A ( $\eta = (\pm, +, +)$  flat metric)

Let  $F_A = dA + \frac{1}{2}[A \wedge A]$  be the curvature of A, and  $\langle, \rangle$  the Killing form on  $\mathfrak{so}(\eta) \cong \mathbb{R}^3$ 

 $\bullet$  Action for 2+1 gravity with cosmological constant  $\Lambda$ 

$$S_{\Lambda} = \int_{M} \langle e \wedge F_{A} \rangle - rac{\Lambda}{3} \langle e \wedge [e \wedge e] 
angle$$
 (1)

• Goal : make sense of the formal quantity

$$\mathcal{Z}(M) = \int D[e]D[A] \exp iS_{\Lambda}[e, A]$$
(2)

• Ponzano-Regge (PR) model : regularisation of  $\mathcal{Z}(M)$  for Euclidean gravity with  $\Lambda = 0$ 

The Ponzano-Regge model

#### Ponzano-Regge model [Ponzano, Regge '68; Freidel, Louapre '04; Barrett, Naish-Guzman '08]

- Let M be a closed, oriented triangulated 3-manifold
- PR model based on Rep(SU(2)). Data consists of :
  - Assignments  $I : e \mapsto I(e) \in \text{Irrep}(SU(2))$  to the edges e of M
  - A state space for each triangle  $\Delta$  of M

$$\mathcal{H}_{\Delta} = \operatorname{Hom}(\bigotimes_{e \in \partial \Delta} I(e), \mathbb{C}) \ni \alpha,$$

• An amplitude  $A_t : \bigotimes_{\Lambda \in \partial t} H_{\Delta} \to \mathbb{C}$  for each tetrahedron t

$$A_t(\alpha \otimes \beta \otimes \gamma \otimes \delta) = \{6j\}_t$$

- $PR(M) = \sum \prod \dim I(e) \prod \{6j\}_t$ PR model : (3)
- Problem : the above infinite sum generally diverges

[Bonzom, Smerlak '10]

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# The Turaev-Viro model [Turaev, Viro '92]

- Goal : regularise the divergencies of the PR model
- Idea : the quantum group  $U_q(\mathfrak{su}(2))$  with  $q^r = 1$  admits only a finite number of irreducible representations [Arnaudon, Roche '89]
  - $\rightarrow$  Natural regularisation of the PR model : replace SU(2) by its quantum deformation  $U_q(\mathfrak{su}(2))$  at root of unity
- Irreducible representations of  $U_q(\mathfrak{su}(2))$  (after purification):

$$I \in \left\{0, rac{1}{2}, 1, ..., rac{r-2}{2}
ight\}, \quad q = \exp 2i\pi/r$$

Constructing a model based on Rep(Uq(su(2))) leads to a finite model: the Turaev-Viro invariant

$$TV_q(M) = K \sum_{I} \prod_{e} [\dim I(e)]_q \prod_t \{6j_q\}_t$$
(4)

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### Relation to the cosmological constant $\Lambda$

• Using the equivalence  $_{[Witten `88]}$  between 2+1 gravity and Chern-Simons theory one can show that :

$$TV_q(M) \propto \int D[e]D[A] \exp iS_{\Lambda}[e,A], \quad \text{if } q = \exp il_p/l_c$$
 (5)

• Asymptotics of the quantum 6j symbol when  $I(e) \rightarrow \infty$  : [Mizoguchi,Tada '92]

$$\{6j_q\}_t \sim \frac{1}{\sqrt{12\pi V_t}} \cos\left(S_t + \frac{\pi}{4}\right), \quad V_t = \operatorname{volume}(t), \quad (6)$$

where  $S_t$  is the Regge action with cosmological constant  $\Lambda$  for the tetrahedron t

$$S_t = \sum_{e \in \partial t} \theta_e I(e) - \frac{1}{l_c^2} V_t$$
(7)

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BF theory The Ooguri model The Crane-Yetter model

4d BF theory [Horowitz '89; Cattaneo, Cotta-Ramusino, Fröhlich, Martellini '95; Baez '95]

• BF theory with semi-simple Lie group G on a 4-manifold M:

- B field  $B \in \Omega^2(M) \otimes \mathfrak{g}$
- connection A on a principal G-bundle over M

Let  $F_A = dA + \frac{1}{2}[A \wedge A]$  be the curvature of A, and  $\langle, \rangle$  the Killing form on  $\mathfrak{g}$ 

 $\bullet$  Action of 4d BF theory with cosmological constant  $\Lambda$ 

$$S_{\Lambda} = \int_{\mathcal{M}} \langle B \wedge F_A \rangle - \frac{\Lambda}{12} \langle B \wedge B \rangle$$
 (8)

• Goal : make sense of the formal quantity

$$\mathcal{Z}(M) = \int D[B]D[A] \exp iS_{\Lambda}[B, A]$$
(9)

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BF theory **The Ooguri model** The Crane-Yetter model

The Ooguri model [Ooguri '92]

- Ooguri model : regularisation of  $\mathcal{Z}(M)$  in the case of  $G = \mathrm{SU}(2)$  and  $\Lambda = 0$
- Model based on Rep(SU(2))
  - The assignments are now to the triangles  $\Delta$  and tetrahedra t of a triangulation of M

$$I: \Delta \mapsto I(\Delta) \in \mathsf{Irrep}(\mathrm{SU}(2)), \quad t \mapsto \alpha_t \in H_t = \mathrm{Hom}(\bigotimes_{\Delta \in \partial t} I(\Delta), \mathbb{C})$$

- The amplitude for the 4-simplexes  $\sigma$  are given by 15*j*-symbols
- Partition function :

$$O(M) = \sum_{I,J} \prod_{\Delta} \dim I(\Delta) \prod_{t} \dim J(t)^{-1} \prod_{\sigma} \{15j\}_{\sigma}$$
(10)

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BF theory The Ooguri model The Crane-Yetter model

The Crane-Yetter model [Crane, Yetter, Kauffman '93]

- The Ooguri model diverges in general
- It can be regularised by considering a model based on  $\operatorname{Rep}(U_q(\mathfrak{su}(2)))$  with  $q^r = 1$

$$CY_q(M) = K \sum_{I,J} \prod_{\Delta} [\dim I(\Delta)]_q \prod_t [\dim J(t)]_q^{-1} \prod_{\sigma} \{15j_q\}_{\sigma}$$
(11)

- The Crane-Yetter model is finite and provides an invariant of topological 4-manifolds
- It is related to SU(2) BF theory on M with cosmological constant  $\Lambda$  via Chern-Simons theory on  $\partial M$  [Roberts '93; Baez '95]  $CY_q(M) \propto \int D[B]D[A] \exp iS_{\Lambda}[B, A], \quad \text{if } q = \exp il_p^2/l_c^2$  (12)

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Gravity as a constrained BF theory [Plebanski '77; Freidel, De Pietri '98]

- First order gravity on a 4-dimensional manifold *M*:
  - dual co-frame  $e: TM \to \mathbb{R}^4$
  - $\mathrm{SO}(\eta)$  connection A ( $\eta=(\epsilon,+,+,+)$ ,  $\epsilon=\pm 1$ , flat metric)
- $\bullet$  Action for 3+1 gravity with cosmological constant  $\Lambda$

$$S_{\Lambda} = \int_{\mathcal{M}} \langle *(e \wedge e) \wedge F_{\mathcal{A}} \rangle - \frac{\Lambda}{12} \langle *(e \wedge e) \wedge e \wedge e \rangle \qquad (13)$$

• Suggests that GR could be cast as a BF theory :

$$S_{\Lambda} = \int_{M} \langle B[e] \wedge F_{A} 
angle - rac{\epsilon \Lambda}{12} \langle B[e] \wedge *B[e] 
angle, \quad ext{with} \quad B[e] = *(e \wedge e)$$

• Gravity  $\equiv$  BF theory + constraints on the *B* field

$$S_{\Lambda}^{\text{Plebanski}} = \int_{M} \langle B \wedge F_{A} \rangle - \frac{\epsilon \Lambda}{12} \langle B \wedge *B \rangle + \mathcal{C}(B)$$
(14)

# The EPRL/FK models [Engle, Pereira, Rovelli, Livine '08; Freidel, Krasnov '08]

- EPRL/FK models : spin foam models for Plebanski's theory with  $\Lambda=$  0 and Immirzi parameter  $\gamma$
- Idea : implement the Plebanski constraints as a choice of measure in the path integral for BF theory
- Concretely :
  - Consider a generalised Ooguri model for  $Rep(SO(\eta))$
  - Impose constraints on the data of the model, i.e., at the level of representation labels  $I(\Delta)$  and state spaces  $H_t$
- Questions :
  - Are the EPRL/FK models finite ? [Perini, Rovelli, Speziale '09]
  - How to introduce a cosmological constant ?

[Bianchi, Krajewski, Rovelli, Vidotto '11]

 Possible answer : consider the quantum deformation of the models [Smolin '95; Major, Smolin '96; Smolin '02; Rovelli '10]

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# Quantum deformation of the EPRL model

- Our work : construction and analysis of a *q*-deformation of the Euclidean and Lorentzian versions of the EPRL model
  - The Euclidean model is based on

 $U_q(\mathfrak{spin}(4)) = U_q(\mathfrak{su}(2)) \otimes U_q(\mathfrak{su}(2)), \quad \text{with $q$ root of unity}$ 

• The Lorentzian model on

 $U_q(\mathfrak{sl}(2,\mathbb{C})) = \mathcal{D}U_q(\mathfrak{su}(2)) = U_q(\mathfrak{su}(2)) \hat{\otimes} F_q(\mathrm{SU}(2))^{op}$ , with q real

- Basic ingredients for the construction of the model :
  - A given subset of representations *I* ⊂ Rep(*U<sub>q</sub>*(so(η))) (*q*-analogue of EPRL representations)
  - A specific class of intertwiners *i* between the elements of *I* (*q*-analogue of EPRL intertwiners)
- From here on, we will focus on the Lorentzian model

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# q-Lorentzian EPRL representations

• Unitary, irreducible representations (of the principal series) of  $U_q(\mathfrak{sl}(2,\mathbb{C}))$ ,  $q = e^{-\kappa}$  [Pusz '93, Buffenoir, Roche '99] :

 $\mathsf{Irrep}(U_q(\mathfrak{sl}(2,\mathbb{C}))) = \{(n,p), n \in \mathbb{Z}/2, p \in [0, 4\pi/\kappa[\} (15) \}$ 

• Let  $\gamma \in \mathbb{R}$  be a parameter. An EPRL representation is a map

$$\phi_{\gamma}: \operatorname{Irrep}(U_q(\mathfrak{sl}(2)) \to \operatorname{Irrep}(U_q(\mathfrak{sl}(2,\mathbb{C})))),$$

defined by

$$K \mapsto (n(K), p(K)) = (K, \gamma K)$$
(16)

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• Remark: the pre-image of the map  $\phi_{\gamma}$  is restricted to

$$\mathcal{L} = \{ \mathcal{K} \in \mathbb{N}/2 \ | \ \mathcal{K} < 4\pi/\gamma\kappa \} \subset \mathsf{Irrep}(U_q(\mathfrak{su}(2)))$$

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### q-Lorentzian EPRL interwiner : Definition I.

- Notation : principal representations ( $\pi_{lpha}, V_{lpha}$ ), lpha = (n, p)
- The EPRL representation  $\alpha(K)$  factorises as

$$V_{\alpha(K)} = \bigoplus_{I=K}^{\infty} V_I, \qquad V_I : U_q(\mathfrak{su}(2)) \text{-module}$$
(17)

• Let  $f_{\alpha}^{\mathcal{K}}: V_{\alpha} \to V_{\mathcal{K}}$  be the projection on the lowest weight factor. The dual map induces an embedding

$$f^*: \operatorname{Hom}_{U_q(\mathfrak{su}(2))}(\otimes_{i=1}^n V_{K_i}, \mathbb{C}) \to \operatorname{Hom}_{U_q(\mathfrak{sl}(2,\mathbb{C}))}(\otimes_{i=1}^n V_{\alpha_i}, \mathbb{C})$$

To all U<sub>q</sub>(su(2))-intertwiner Λ<sub>K</sub> : ⊗<sup>n</sup><sub>i=1</sub>V<sub>Ki</sub> → C, this map associates a quantum EPRL intertwiner ι<sub>α(K)</sub>

$$\iota_{\alpha(K)} = f^*(\Lambda_K) = \Lambda_K \circ \bigotimes_{i=1}^n f_{\alpha_i}^{K_i} \circ T_{\alpha_1,\dots,\alpha_n}$$
(18)

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# *q*-Lorentzian EPRL interwiner : Definition II.

•  $T_{\alpha_1,...,\alpha_n}$  is the *q*-analogue of the classical expression

$$\int_{\mathrm{SL}(2,\mathbb{C})} dX \left(\bigotimes_{i=1}^n \pi_{\alpha_i}\right) (X)$$
(19)

• Let  $h: F_q(\mathrm{SL}(2,\mathbb{C})) \to \mathbb{C}$  be a Haar measure on the Hopf algebra  $F_q(\mathrm{SL}(2,\mathbb{C}))$  dual to  $U_q(\mathfrak{sl}(2,\mathbb{C}))$  [Buffenoir, Roche '99] Ex. :  $f \in F(\mathrm{SL}(2,\mathbb{C})), \quad h(f) = \int_{\mathrm{SL}(2,\mathbb{C})} dX \ f(X)$ 

•  $\mathcal{T}_{\alpha_1,...,\alpha_n}$  is defined as [Noui, Roche '02]

$$T_{\alpha_1,\ldots,\alpha_n} = \sum_{\mathcal{A}} \left( \bigotimes_{i=1}^n \pi_{\alpha_i} \right) \left( \Delta^{(n)}(x^{\mathcal{A}}) \right) h(x_{\mathcal{A}}), \qquad (20)$$

where  $\{x^A\}_A$  is a basis of  $U_q(\mathfrak{sl}(2,\mathbb{C}))$  and  $\{x_A\}_A$  is the dual basis of  $F_q(\mathrm{SL}(2,\mathbb{C}))$ 

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# q-Lorentzian EPRL interwiner : Properties

- Theorem. 1. Let {e<sup>I</sup><sub>a</sub>(α) | I ∈ N, I ≥ |n|, a = -I, ..., I} be a basis of V<sub>α</sub>. The evaluation of the 4-valent quantum EPRL intertwiner ι<sub>α</sub>(⊗<sup>4</sup><sub>i=1</sub> e<sup>I<sub>i</sub></sup><sub>a<sub>i</sub></sub>(α<sub>i</sub>)) is a multiple series which converges absolutely
  2. The linear map ι<sub>α</sub> : ⊗<sup>4</sup><sub>i=1</sub>V<sub>α<sub>i</sub></sub> → C is a intertwiner for the quantum Lorentz group
- **Proposition.** Let *R* be the matrix of the quantum double  $\mathcal{D}U_q(\mathfrak{su}(2))$  and  $c_{\alpha_2,\alpha_1}: V_{\alpha_2} \otimes V_{\alpha_1} \to V_{\alpha_1} \otimes V_{\alpha_2}$  with

$$c_{\alpha_2,\alpha_1} = \tau_{\alpha_2,\alpha_1} \circ (\pi_{\alpha_2} \otimes \pi_{\alpha_1})(R), \qquad (21)$$

the associated braiding. The quantum EPRL intertwiner  $\iota_{\alpha}$  is not invariant under the action of c :

$$\iota_{\alpha_{1}\alpha_{2}\alpha_{3}\alpha_{4}} \circ c_{\alpha_{2},\alpha_{1}} \neq \iota_{\alpha_{2}\alpha_{1}\alpha_{3}\alpha_{4}}$$
(22)

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# Graphical calculus : ingredients

- 4-simplex amplitude defined using graphical calculus
- Elements for the diagrammatic calculus :

• EPRL tensors : elements of 
$$\left[\bigotimes_{i=1}^{4} V_{\alpha_{i}}\right] \otimes F_{q}(\mathrm{SL}(2,\mathbb{C}))$$
  
 $\Psi^{\alpha} = \left[\sum_{A} \left(\bigotimes_{i=1}^{4} \pi_{\alpha_{i(K_{i})}} \left(\Delta^{(4)}(x^{A})\right)\right) \circ \bigotimes_{i=1}^{4} f_{K_{i}}^{\alpha_{i}} \circ \Lambda^{K}\right] \otimes x_{A} :=$ 
where  $\Lambda^{K} \in \mathrm{Hom}_{U_{q}(\mathfrak{su}(2))} \left(\mathbb{C}, \bigotimes_{i=1}^{4} V_{K_{i}}\right) \text{ and } f_{K}^{\alpha} : V_{K} \to V_{\alpha}$   
• An invariant bilinear form  $\beta : V_{\alpha} \times V_{\alpha} \to \mathbb{C}$   
 $\beta :=$ 
• The braiding  $c_{\alpha,\beta} : V_{\alpha} \otimes V_{\beta} \to V_{\beta} \otimes V_{\alpha}$   
 $c_{\alpha,\beta} :=$ 

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# Graphical calculus : evaluation

- Closed diagram  $\Gamma_n$  with *n* vertices  $\rightarrow \phi(\Gamma_n) \in F_q(\mathrm{SL}(2,\mathbb{C}))^{\otimes n}$
- The evaluation of a closed diagram is defined as

$$ev(\Gamma_n) = (\epsilon \otimes h^{n-1})(\phi(\Gamma_n)), \quad \epsilon \text{ is the co-unit}$$
 (23)

• The diagram for the 4-simplex amplitude is given by



(24)

• **Theorem.** The multiple series  $ev(\Gamma_5)$  converges absolutely

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# The model

- Let M be a closed, oriented triangulated 4-manifold
- Model based on  $\operatorname{Rep}(U_q(\mathfrak{sl}(2,\mathbb{C})))$ . Data consists of :
  - An assignment  $\alpha : \Delta \mapsto \alpha(\Delta) \in \operatorname{Irrep}(U_q(\mathfrak{sl}(2,\mathbb{C})))$  of an EPRL representation to each triangle  $\Delta$  of M
  - A state space of EPRL tensors for each tetrahedron t of M

$$H_t = \left(\bigotimes_{\Delta \in \partial t} \alpha(\Delta)\right) \otimes F_q(\mathrm{SL}(2,\mathbb{C}))$$

• An amplitude  $A_{\sigma}: \bigotimes_{t \in \partial \sigma} H_t \to \mathbb{C}$  for each 4-simplex  $\sigma$ 

$$A_{\sigma}(\Psi_1\otimes...\otimes\Psi_5)=ev({\sf \Gamma}_5)\ \in\mathbb{C}$$

• The resulting model is *finite*. It is given by

$$\mathcal{Z}_{q}(M) = \sum_{K \in \mathcal{L}} \sum_{J} \prod_{\Delta} [\dim K(\Delta)]_{q} \prod_{\sigma} A_{\sigma}(\{K(\Delta)\}, \{J(t)\})$$
(25)

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# Conclusion

- We have constructed and analysed a q-deformation of the EPRL model. We have :
  - Generalised the classical constructions to models based on  $U_q(\mathfrak{su}(2)) \otimes U_q(\mathfrak{su}(2))$  and  $\mathcal{D}U_q(\mathfrak{su}(2))$
  - Defined the EPRL intertwiner, studied its convergence and properties under braiding
  - Constructed a convergent amplitude for the 4-simplexes
- Open question: relation to the cosmological constant  $\Lambda$  ?

• If  $q = e^{-l_p^2/l_c^2}$  (Lorentzian), bound on the area of the triangles

$$A(\Delta) < 32\pi^2 l_c^2$$
  $(l_p << l_c)$  (26)

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• Need for an asymptotic formula [Ding, Han '11]