

# Some analytical results of the Hamiltonian operator in LQG

Cong Zhang

January 22, 2018





# Outline

Deparametrized model of LQG coupled to scalar field

General works about the physical Hamiltonian

Restriction on a special case

A toy model of LQG cosmology

Conclusion and outlook



# Deparametrized model of LQG coupling to a scalar field

(Alesci et al., 2015)

- Introduce the other fields as reference frame classically.

# Deparametrized model of LQG coupling to a scalar field

(Alesci et al., 2015)

- ▶ Introduce the other fields as reference frame classically.
- ▶ Solve the constraint equation to obtain a physical Hamiltonian. In the model coupling to a scalar field

$$C'(x) = \pi(x) \pm \sqrt{h(x)}$$

where

$$h(x) = -\sqrt{|\det E|} \left( -C^{\text{gr}} \pm \sqrt{(C^{\text{gr}})^2 - q^{ab} C_a^{\text{gr}} C_b^{\text{gr}}} \right)$$

# Deparametrized model of LQG coupling to a scalar field

(Alesci et al., 2015)

- ▶ Introduce the other fields as reference frame classically.
- ▶ Solve the constraint equation to obtain a physical Hamiltonian. In the model coupling to a scalar field

$$C'(x) = \pi(x) \pm \sqrt{h(x)}$$

where

$$h(x) = -\sqrt{|\det E|} \left( -C^{\text{gr}} \pm \sqrt{(C^{\text{gr}})^2 - q^{ab} C_a^{\text{gr}} C_b^{\text{gr}}} \right)$$

- ▶ Quantize this system.

# Main results of the quantum Theory (Domagała et al., 2010,

Lewandowski et al., 2011)

- **The physical Hilbert space  $\mathcal{H}_{\text{phy}}$ :**  
the Hilbert space space of pure gravity, satisfying Gaussian and vector constraints.

# Main results of the quantum Theory (Domagała et al., 2010, Lewandowski et al., 2011)

- ▶ **The physical Hilbert space  $\mathcal{H}_{\text{phy}}$ :**  
the Hilbert space space of pure gravity, satisfying Gaussian and vector constraints.
- ▶ **The dynamic:**

$$i\hbar \frac{d}{dt} \Psi = \hat{H} \Psi$$

where  $t$  is a parameter of the transformations  $\phi \mapsto \phi + t$ .

# Main results of the quantum Theory<sup>(Domagała et al., 2010,</sup>

Lewandowski et al., 2011)

- ▶ **The physical Hilbert space**  $\mathcal{H}_{\text{phy}}$ :  
the Hilbert space space of pure gravity, satisfying Gaussian and vector constraints.
- ▶ **The dynamic:**

$$i\hbar \frac{d}{dt} \Psi = \hat{H} \Psi$$

where  $t$  is a parameter of the transformations  $\phi \mapsto \phi + t$ .

- ▶ **The quantum Hamiltonian**

$$\hat{H} = \int d^3x \sqrt{-2 \sqrt{\overline{|\det E(x)|}} C^{\text{gr}}(x)}$$



# Main results of the quantum Theory (Domagała et al., 2010,

Lewandowski et al., 2011)

- ▶ **The physical Hilbert space**  $\mathcal{H}_{\text{phy}}$ :  
the Hilbert space space of pure gravity, satisfying Gaussian and vector constraints.
- ▶ **The dynamic:**

$$i\hbar \frac{d}{dt} \Psi = \hat{H} \Psi$$

where  $t$  is a parameter of the transformations  $\phi \mapsto \phi + t$ .

- ▶ **The quantum Hamiltonian**

$$\hat{H} = \int d^3x \sqrt{-2 \sqrt{\overline{|\det E(x)|}} C^{\text{gr}}(x)}$$

What I do is to study this (physical) Hamiltonian operator.

## General works about the physical Hamiltonian

- Classical expression of  $\sqrt{|\det E(x)|} C^{\text{gr}}(x)$

$$\begin{aligned} & - \sqrt{|\det E(x)|} C^{\text{gr}}(x) \\ &= \frac{1}{16\pi\beta^2 G} (\epsilon_{ijk} E_i^a(x) E_j^b(x) F_{ab}^k(x) + (1 + \beta^2) |\det E(x)| R(x)) \\ &=: \frac{1}{16\pi\beta^2 G} (H^E(x) + H^L(x)). \end{aligned}$$

# General works about the physical Hamiltonian

- Classical expression of  $\sqrt{|\det E(x)|} C^{\text{gr}}(x)$

$$\begin{aligned} & - \sqrt{|\det E(x)|} C^{\text{gr}}(x) \\ &= \frac{1}{16\pi\beta^2 G} (\epsilon_{ijk} E_i^a(x) E_j^b(x) F_{ab}^k(x) + (1 + \beta^2) |\det E(x)| R(x)) \\ &=: \frac{1}{16\pi\beta^2 G} (H^E(x) + H^L(x)). \end{aligned}$$

- Volume operator is not involved.

# General works about the physical Hamiltonian

- ▶ Classical expression of  $\sqrt{|\det E(x)|} C^{\text{gr}}(x)$

$$\begin{aligned} & - \sqrt{|\det E(x)|} C^{\text{gr}}(x) \\ &= \frac{1}{16\pi\beta^2 G} (\epsilon_{ijk} E_i^a(x) E_j^b(x) F_{ab}^k(x) + (1 + \beta^2) |\det E(x)| R(x)) \\ &=: \frac{1}{16\pi\beta^2 G} (H^E(x) + H^L(x)). \end{aligned}$$

- ▶ Volume operator is not involved.
  - ▶ The analysis is simpler.

## General works about the physical Hamiltonian

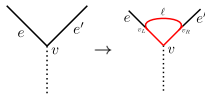
- ▶ Classical expression of  $\sqrt{|\det E(x)|} C^{\text{gr}}(x)$

$$\begin{aligned} & - \sqrt{|\det E(x)|} C^{\text{gr}}(x) \\ &= \frac{1}{16\pi\beta^2 G} (\epsilon_{ijk} E_i^a(x) E_j^b(x) F_{ab}^k(x) + (1 + \beta^2) |\det E(x)| R(x)) \\ &=: \frac{1}{16\pi\beta^2 G} (H^E(x) + H^L(x)). \end{aligned}$$

- ▶ Volume operator is not involved.
  - ▶ The analysis is simpler.
  - ▶ It is possible to start from the simplest case of 2-valent graph.

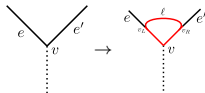
# General works about the physical Hamiltonian

- The chosen loop to quantize  $F_{ab}$  (Thiemann, 2007, Yang and Ma, 2015)



# General works about the physical Hamiltonian

- The chosen loop to quantize  $F_{ab}$  (Thiemann, 2007, Yang and Ma, 2015)



- Expression of the Hamiltonian operator (Alesci et al., 2015)

$$\hat{H} = \frac{1}{\sqrt{16\pi G\beta^2}} \sum_{v \in V(\gamma)} \sqrt{\sum_{e, e' \text{ at } v} \epsilon(e, e') (H_{v, ee'}^E + (H_{v, ee'}^E)^\dagger) + (1 + \beta^2) H_{v, ee'}^L}$$

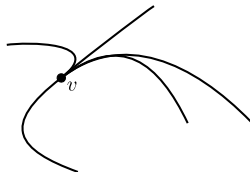
$$H_{v, ee'}^E = \epsilon_{ijk} \text{Tr}^{(I)}(h_{\alpha_{ee'}}^i \tau^j) J_{v, e}^j J_{v, e'}^k$$

$$H_{ee'}^L := \sqrt{\delta_{ii'} (\epsilon_{ijk} J_{v, e}^j J_{v, e'}^k) (\epsilon_{i'j'k'} J_{v, e}^{j'} J_{v, e'}^{k'})} \left( \frac{2\pi}{\alpha} - \pi + \arccos \left[ \frac{\delta_{kl} J_{v, e}^k J_{v, e'}^l}{\sqrt{\delta_{kk'} J_{v, e}^k J_{v, e'}^{k'}} \sqrt{\delta_{kk'} J_{v, e'}^k J_{v, e}^{k'}}} \right] \right)$$

# General works about the physical Hamiltonian

The physical Hilbert space:

- Degenerate vertex:

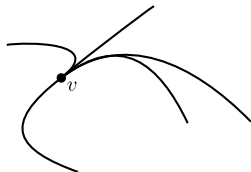




# General works about the physical Hamiltonian

The physical Hilbert space:

- Degenerate vertex:

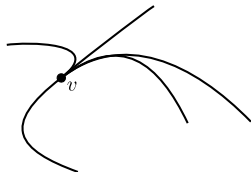


- $V_{nd}(\gamma)$ : non-degenerate vertices.  
 $\text{Diff}_{V_{nd}}$ : diffeomorphisms preserving  $V_{nd}$ .  
 $\text{Diff}(\gamma)_{\text{Tr}}$ : diffeomorphisms acting trivially on  $\gamma$

# General works about the physical Hamiltonian

The physical Hilbert space:

- Degenerate vertex:

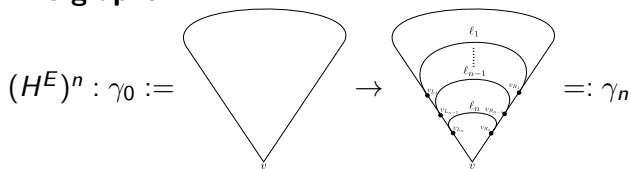


- $V_{\text{nd}}(\gamma)$ : non-degenerate vertices.  
 $\text{Diff}_{V_{\text{nd}}}$ : diffeomorphisms preserving  $V_{\text{nd}}$ .  
 $\text{Diff}(\gamma)_{\text{Tr}}$ : diffeomorphisms acting trivially on  $\gamma$
- Physical state from  $|\psi_\gamma\rangle$

$$(\psi_\gamma| = \mathcal{N}_\gamma \sum_{\phi \in \text{Diff}_{V_{\text{nd}}} / \text{Diff}(\gamma)_{\text{Tr}}} U_\phi |\psi_\gamma\rangle := \eta |\psi\rangle$$

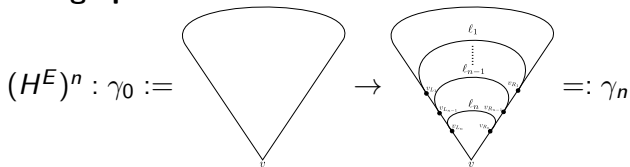
# Restriction on a special case

- The graphs:



## Restriction on a special case

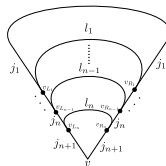
- The graphs:



- The Hilbert space: Spin networks on  $\gamma_n$ .

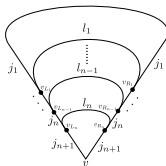
## Restriction on a special case

- Spin networks on  $\gamma_n$ :  $|\gamma_n, \vec{j}, \vec{l}\rangle \sim$



## Restriction on a special case

- Spin networks on  $\gamma_n$ :  $|\gamma_n, \vec{j}, \vec{l}\rangle \sim$

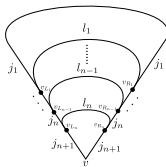


- Ambiguity of choosing spin  $l$ :

$$\epsilon_{ijk} \text{Tr}^{(l)}(h_{\alpha_{ee'}} \tau^i) J_{v,e}^j J_{v,e'}^k : \quad \rightarrow \sum_{j_{n+1}} C_{j_{n+1}} \quad .$$

## Restriction on a special case

- Spin networks on  $\gamma_n$ :  $|\gamma_n, \vec{j}, \vec{l}\rangle \sim$



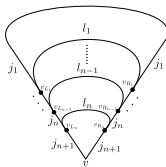
- Ambiguity of choosing spin  $l$ :

$$\epsilon_{ijk} \text{Tr}^{(l)}(h_{\alpha_{ee'}} \tau^i) J_{v,e}^j J_{v,e'}^k : \quad \rightarrow \sum_{j_{n+1}} C_{j_{n+1}} \quad .$$

- $l(j_n) = \begin{cases} 1 & , j_n = 1/2, \\ 1/2 & , j_n \neq 1/2. \end{cases}$

## Restriction on a special case

- Spin networks on  $\gamma_n$ :  $|\gamma_n, \vec{j}, \vec{l}\rangle \sim$



- Ambiguity of choosing spin  $l$ :

$$\epsilon_{ijk} \text{Tr}^{(l)}(h_{\alpha_{ee'}} \tau^i) J_{v,e}^j J_{v,e'}^k : \quad \rightarrow \quad \sum_{j_{n+1}} C_{j_{n+1}} \quad \cdot$$

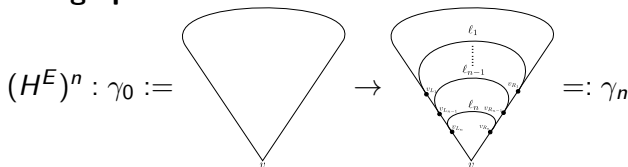
- $l(j_n) = \begin{cases} 1 & , j_n = 1/2, \\ 1/2 & , j_n \neq 1/2. \end{cases}$

- $\mathcal{H}_{\text{phy}} = \overline{\text{span}([ \gamma_n ], \vec{j} \mid := \mathcal{N} \sum_{\phi \in \text{Diff}_v / \text{Diff}_{\text{Tr}}(\gamma_n)} U_\phi | \gamma_n, \vec{j}, l(\vec{j}) \rangle)}$



## Restriction on a special case

- The graphs:

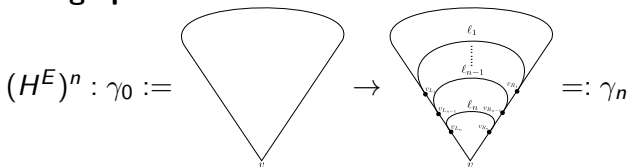


- Hilbert space:

$$\mathcal{H}_{\text{phy}} = \text{span}([ \gamma_n ], \vec{j} \mid := \mathcal{N} \sum_{\phi \in \text{Diff}_v / \text{Diff}_{\text{Tr}}(\gamma_n)} U_\phi | \gamma_n, \vec{j}, l(\vec{j}) \rangle).$$

## Restriction on a special case

- **The graphs:**

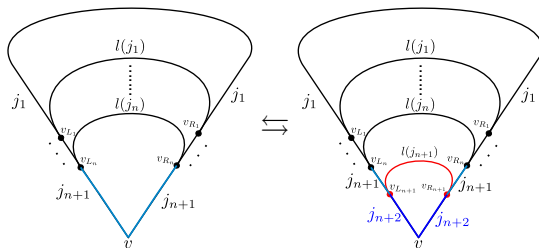


- **Hilbert space:**

$$\mathcal{H}_{\text{phy}} = \text{span}([ \gamma_n ], \vec{j} \mid := \mathcal{N} \sum_{\phi \in \text{Diff}_V / \text{Diff}_{\text{Tr}}(\gamma_n)} U_\phi | \gamma_n, \vec{j}, l(\vec{j}) \rangle).$$

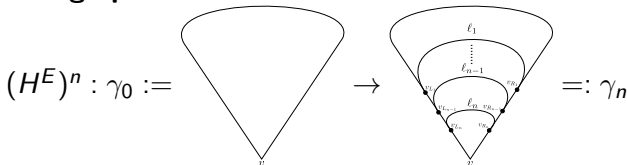
- **Operator:** Hamiltonian operator restricted on the Hilbert space.

## Restriction on a special case



## Restriction on a special case

- The graphs:



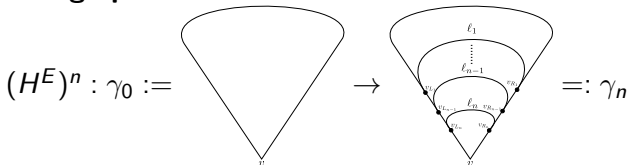
- Hilbert space:

$$\mathcal{H}_{\text{phy}} = \text{span}([[\gamma_n], \vec{j} \mid := \mathcal{N} \sum_{\phi \in \text{Diff}_V / \text{Diff}_{\text{Tr}}(\gamma_n)} U_\phi | \gamma_n, \vec{j}, l(\vec{j}) \rangle).$$

- Operator: Hamiltonian operator restricted on the Hilbert space,  $H|_{\mathcal{H}_{\text{phy}}}$ .

## Restriction on a special case

- The graphs:



- Hilbert space:

$$\mathcal{H}_{\text{phy}} = \text{span}([[\gamma_n], \vec{j} \mid := \mathcal{N} \sum_{\phi \in \text{Diff}_V / \text{Diff}_{\text{Tr}}(\gamma_n)} U_\phi | \gamma_n, \vec{j}, l(\vec{j}) \rangle).$$

- **Operator:** Hamiltonian operator restricted on the Hilbert space,  $H|_{\mathcal{H}_{\text{phy}}}$ .
- **Question:** Self-adjointness of the restricted  $H$ .

# Self-adjointness of the Hamiltonian operator

## Theorem

Let  $N$  be a self-adjoint operator with  $N \geq 1$ . Let  $H$  be a symmetric operator with domain  $D$  which is a core for  $N$ . Suppose that:

i For some  $c$  and all  $\psi \in D$ ,

$$\|H\psi\| \leq c\|N\psi\|.$$

ii For some  $d$  and all  $\psi \in D$ ,

$$|(H\psi, N\psi) - (N\psi, H\psi)| \leq d\|N^{1/2}\psi\|^2$$

Then  $A$  is essential self-adjoint on  $D$  and its closure is essentially self-adjoint on any core for  $N$ .

# Self-adjointness of the Hamiltonian operator

## Theorem

Let  $N$  be a self-adjoint operator with  $N \geq 1$ . Let  $H$  be a symmetric operator with domain  $D$  which is a core for  $N$ . Suppose that:

i For some  $c$  and all  $\psi \in D$ ,

$$\|H\psi\| \leq c\|N\psi\|.$$

ii For some  $d$  and all  $\psi \in D$ ,

$$|(H\psi, N\psi) - (N\psi, H\psi)| \leq d\|N^{1/2}\psi\|^2$$

Then  $A$  is essential self-adjoint on  $D$  and its closure is essentially self-adjoint on any core for  $N$ .

We can choose the operator  $N$ , diagonalized under the basis, that is  $([\gamma_n], \vec{j} | N = ([\gamma_n], \vec{j} | N(j_{n+1}))$ , such that  $N(j) \cong j^n$  with  $n \geq 1$ .



## Summation

- We prove the operator  $H$ , restricted on the simplest graph, is self-adjoint.





## Summation

- ▶ We prove the operator  $H$ , restricted on the simplest graph, is self-adjoint.
- ▶ Turn to the graph preserving version.

# Summation

- ▶ We prove the operator  $H$ , restricted on the simplest graph, is self-adjoint.
- ▶ Turn to the graph preserving version.
  - ▶ Consider the coherent state peaking at the cosmology phase space.

# Summation

- ▶ We prove the operator  $H$ , restricted on the simplest graph, is self-adjoint.
- ▶ Turn to the graph preserving version.
  - ▶ Consider the coherent state peaking at the cosmology phase space.
  - ▶ Calculate the semiclassical dynamics with the coherent state as the initial data.

# The heat kernel coherent state

- Classical phase space of the cosmology phase space:

$$A_a^i = cV_o^{-1/3}\dot{\omega}_a^i, E_i^a = pV_o^{-2/3}\sqrt{q_o}\dot{e}_i^a$$

# The heat kernel coherent state

- ▶ Classical phase space of the cosmology phase space:

$$A_a^i = cV_o^{-1/3}\dot{\omega}_a^i, E_i^a = pV_o^{-2/3}\sqrt{q_o}\dot{e}_i^a$$

- ▶ The graph: dipole graph with two  $N$ -valent vertices.

# The heat kernel coherent state

- ▶ Classical phase space of the cosmology phase space:

$$A_a^i = cV_o^{-1/3}\dot{\omega}_a^i, E_i^a = pV_o^{-2/3}\sqrt{q_o}\dot{e}_i^a$$

- ▶ The graph: dipole graph with two  $N$ -valent vertices.
- ▶ Coherent state for  $p \gg 1$  (Bahr and Thiemann, 2009, Bianchi et al., 2010):

$$\Psi_{c,p}(\vec{g}) \cong \prod_{e \in E(\gamma)} \left( \sum_{j_e} (2j_e + 1) e^{-tj_e(j_e+1) + \nu pj_e - i\mu c j_e} D_{j_e, j_e}^{j_e} (n_e^{-1} g_e n_e) \right)$$

# The heat kernel coherent state

- ▶ Classical phase space of the cosmology phase space:

$$A_a^i = cV_o^{-1/3}\dot{\omega}_a^i, E_i^a = pV_0^{-2/3}\sqrt{q_o}\dot{e}_i^a$$

- ▶ The graph: dipole graph with two  $N$ -valent vertices.
- ▶ Coherent state for  $p \gg 1$  (Bahr and Thiemann, 2009, Bianchi et al., 2010):

$$\Psi_{c,p}(\vec{g}) \cong \prod_{e \in E(\gamma)} \left( \sum_{j_e} (2j_e + 1) e^{-tj_e(j_e+1) + \nu pj_e - i\mu cj_e} D_{j_e, j_e}^{j_e} (n_e^{-1} g_e n_e) \right)$$

- ▶ The coherent states inspires us to consider the Hilbert space

$$\mathcal{H}_{\text{cos}} = \overline{\text{span}(\langle \vec{g} | \vec{j} \rangle := \bigotimes_{e \in \gamma} D_{j_e, j_e}^{j_e} (n_e^{-1} g n_e))}$$

# The heat kernel coherent state

- ▶ Classical phase space of the cosmology phase space:

$$A_a^i = cV_o^{-1/3}\dot{\omega}_a^i, E_i^a = pV_o^{-2/3}\sqrt{q_o}\dot{e}_i^a$$

- ▶ The graph: dipole graph with two  $N$ -valent vertices.
- ▶ Coherent state for  $p \gg 1$  (Bahr and Thiemann, 2009, Bianchi et al., 2010):

$$\Psi_{c,p}(\vec{g}) \cong \prod_{e \in E(\gamma)} \left( \sum_{j_e} (2j_e + 1) e^{-tj_e(j_e+1) + \nu pj_e - i\mu cj_e} D_{j_e, j_e}^{j_e} (n_e^{-1} g_e n_e) \right)$$

- ▶ The coherent states inspires us to consider the Hilbert space

$$\mathcal{H}_{\text{cos}} = \overline{\text{span}(\langle \vec{g} | \vec{j} \rangle := \bigotimes_{e \in \gamma} D_{j_e, j_e}^{j_e} (n_e^{-1} g_n))}$$

- ▶ The factor  $e^{-tj(j+1) + \nu pj - i\mu cj}$  is a Gaussian function on  $j$  peaking at  $j_0 \cong \frac{\nu p}{2t}$ , so we focus on large  $j$  limit when calculating.



## Action of holonomy and flux operators on $\mathcal{H}_{\text{cos}}$

- For the holonomy operator:

$$h_e^{1/2} D_{n_e n_e}^{j_e}(g) \\ \cong D^{1/2}(n_e) \cdot \begin{pmatrix} D_{n_e n_e}^{j_e+1/2}(g) & 0 \\ 0 & D_{n_e n_e}^{j_e-1/2}(g) \end{pmatrix} \cdot D^{1/2}(n_e^{-1}) + O(1/\sqrt{j})$$

with the abbreviation of  $D_{nn}^j(g) := D_{jj}^j(n^{-1}gn)$

## Action of holonomy and flux operators on $\mathcal{H}_{\text{cos}}$

- For the holonomy operator:

$$h_e^{1/2} D_{n_e n_e}^{j_e}(g) \\ \cong D^{1/2}(n_e) \cdot \begin{pmatrix} D_{n_e n_e}^{j_e+1/2}(g) & 0 \\ 0 & D_{n_e n_e}^{j_e-1/2}(g) \end{pmatrix} \cdot D^{1/2}(n_e^{-1}) + O(1/\sqrt{j})$$

with the abbreviation of  $D_{nn}^j(g) := D_{jj}^j(n^{-1}gn)$

- For the Flux operator

$$\vec{J}_{v,e} D_{n_e n_e}^{j_e}(g) = j_e \vec{n}_e D_{n_e n_e}^{j_e}(g) + O(\sqrt{j}) \quad (1)$$

## Action of holonomy and flux operators on $\mathcal{H}_{\text{cos}}$

- For the holonomy operator:

$$h_e^{1/2} D_{n_e n_e}^{j_e}(g) \\ \cong D^{1/2}(n_e) \cdot \begin{pmatrix} D_{n_e n_e}^{j_e+1/2}(g) & 0 \\ 0 & D_{n_e n_e}^{j_e-1/2}(g) \end{pmatrix} \cdot D^{1/2}(n_e^{-1}) + O(1/\sqrt{j})$$

with the abbreviation of  $D_{nn}^j(g) := D_{jj}^j(n^{-1}gn)$

- For the Flux operator

$$\vec{J}_{v,e} D_{n_e n_e}^{j_e}(g) = j_e \vec{n}_e D_{n_e n_e}^{j_e}(g) + O(\sqrt{j}) \quad (1)$$

- The Hilbert space  $\mathcal{H}_{\text{cos}}$  is preserved approximately for large  $j$ .

# The Hamiltonian operator

► Action of  $H$  on  $\mathcal{H}_{\text{cos}}$ ,  $H = \sum_v \sqrt{|\sum_{ee'} H_{v,ee'}^E + H_{v,ee'}^L|}$

$$\begin{aligned}
 H_{v,ee'}^{(E)} D_{n_e n_e}^{j_e} (g_e) D_{n_{e'}, n_{e'}}^{j_{e'}} (g_{e'}) &\cong (\sqrt{j_e + 1/2} D_{n_e n_e}^{j_e + 1/2} (g_e) - \sqrt{j_e - 1/2} D_{n_e n_e}^{j_e - 1/2} (g_e)) \sqrt{j_e} \times \\
 &\quad \times (\sqrt{j_{e'} + 1/2} D_{n_{e'}, n_{e'}}^{j_{e'} + 1/2} (g_{e'}) - \sqrt{j_{e'} - 1/2} D_{n_{e'}, n_{e'}}^{j_{e'} - 1/2} (g_{e'})) \sqrt{j_{e'}} \\
 H_{v,ee'}^L D_{n_e n_e}^{j_e} (g_e) D_{n_{e'}, n_{e'}}^{j_{e'}} (g_{e'}) &\cong \alpha_{ee'} \sin(\theta_{ee'}) j_e j_{e'} D_{n_e n_e}^{j_e} (g_e) D_{n_{e'}, n_{e'}}^{j_{e'}} (g_{e'})
 \end{aligned}$$

# The Hamiltonian operator

- Action of  $H$  on  $\mathcal{H}_{\text{cos}}$ ,  $H = \sum_v \sqrt{|\sum_{ee'} H_{v,ee'}^E + H_{v,ee'}^L|}$

$$H_{v,ee'}^{(E)} D_{n_e n_e}^{j_e} (g_e) D_{n_{e'}, n_{e'}}^{j_{e'}} (g_{e'}) \cong (\sqrt{j_e + 1/2} D_{n_e n_e}^{j_e + 1/2} (g_e) - \sqrt{j_e - 1/2} D_{n_e n_e}^{j_e - 1/2} (g_e)) \sqrt{j_e} \times \\ \times (\sqrt{j_{e'} + 1/2} D_{n_{e'}, n_{e'}}^{j_{e'} + 1/2} (g_{e'}) - \sqrt{j_{e'} - 1/2} D_{n_{e'}, n_{e'}}^{j_{e'} - 1/2} (g_{e'})) \sqrt{j_{e'}}$$

$$H_{v,ee'}^L D_{n_e n_e}^{j_e} (g_e) D_{n_{e'}, n_{e'}}^{j_{e'}} (g_{e'}) \cong \alpha_{ee'} \sin(\theta_{ee'}) j_e j_{e'} D_{n_e n_e}^{j_e} (g_e) D_{n_{e'}, n_{e'}}^{j_{e'}} (g_{e'})$$

- $H_{v,ee}^E$  and  $H_{v,ee}^L$  are self-adjoint.

# The Hamiltonian operator

- Action of  $H$  on  $\mathcal{H}_{\text{cos}}$ ,  $H = \sum_v \sqrt{|\sum_{ee'} H_{v,ee'}^E + H_{v,ee'}^L|}$

$$H_{v,ee'}^{(E)} D_{n_e n_e}^{j_e} (g_e) D_{n_{e'} n_{e'}}^{j_{e'}} (g_{e'}) \cong (\sqrt{j_e + 1/2} D_{n_e n_e}^{j_e+1/2} (g_e) - \sqrt{j_e - 1/2} D_{n_e n_e}^{j_e-1/2} (g_e)) \sqrt{j_e} \times \\ \times (\sqrt{j_{e'} + 1/2} D_{n_{e'} n_{e'}}^{j_{e'}+1/2} (g_{e'}) - \sqrt{j_{e'} - 1/2} D_{n_{e'} n_{e'}}^{j_{e'}-1/2} (g_{e'})) \sqrt{j_{e'}}$$

$$H_{v,ee'}^L D_{n_e n_e}^{j_e} (g_e) D_{n_{e'} n_{e'}}^{j_{e'}} (g_{e'}) \cong \alpha_{ee'} \sin(\theta_{ee'}) j_e j_{e'} D_{n_e n_e}^{j_e} (g_e) D_{n_{e'} n_{e'}}^{j_{e'}} (g_{e'})$$

- $H_{v,ee}^E$  and  $H_{v,ee}^L$  are self-adjoint.
- $H_{v,ee}^E$  can be rewritten as  $H_{v,ee'}^E = -H_{v,e} H_{v,e'}$  with

$$H_{v,e} D_{n_e n_e}^{j_e} (g_e) = i \left( \sqrt{j_e + 1/2} D_{n_e n_e}^{j_e+1/2} (g_e) - \sqrt{j_e - 1/2} D_{n_e n_e}^{j_e-1/2} (g_e) \right) \sqrt{j_e}$$

# The Hamiltonian operator

- Action of  $H$  on  $\mathcal{H}_{\text{COS}}$ ,  $H = \sum_v \sqrt{|\sum_{ee'} H_{v,ee'}^E + H_{v,ee'}^L|}$

$$H_{v,ee'}^{(E)} D_{n_e n_e}^{j_e} (g_e) D_{n_{e'}, n_{e'}}^{j_{e'}} (g_{e'}) \cong (\sqrt{j_e + 1/2} D_{n_e n_e}^{j_e+1/2} (g_e) - \sqrt{j_e - 1/2} D_{n_e n_e}^{j_e-1/2} (g_e)) \sqrt{j_e} \times \\ \times (\sqrt{j_{e'} + 1/2} D_{n_{e'}, n_{e'}}^{j_{e'}+1/2} (g_{e'}) - \sqrt{j_{e'} - 1/2} D_{n_{e'}, n_{e'}}^{j_{e'}-1/2} (g_{e'})) \sqrt{j_{e'}}$$

$$H_{v,ee'}^L D_{n_e n_e}^{j_e} (g_e) D_{n_{e'}, n_{e'}}^{j_{e'}} (g_{e'}) \cong \alpha_{ee'} \sin(\theta_{ee'}) j_e j_{e'} D_{n_e n_e}^{j_e} (g_e) D_{n_{e'}, n_{e'}}^{j_{e'}} (g_{e'})$$

- $H_{v,ee}^E$  and  $H_{v,ee}^L$  are self-adjoint.
- $H_{v,ee}^E$  can be rewritten as  $H_{v,ee'}^E = -H_{v,e} H_{v,e'}$  with

$$H_{v,e} D_{n_e n_e}^{j_e} (g_e) = i \left( \sqrt{j_e + 1/2} D_{n_e n_e}^{j_e+1/2} (g_e) - \sqrt{j_e - 1/2} D_{n_e n_e}^{j_e-1/2} (g_e) \right) \sqrt{j_e}$$

- For large  $j$ ,

$$H_{v,e} \cong i \sqrt{j_e} \frac{d}{dj_e} \sqrt{j_e} =: H_e^C$$

which is a self-adjoint in the Hilbert space  $L^2(\mathbb{R}^+)$  with

“eigenvector”  $\varphi_\omega(j_e) = \frac{e^{-i\omega \ln(j_e)}}{\sqrt{j_e}}$ .

## Some issues about the Minkowski condition

- Minkowski condition:

$$\sum_{e \text{ at } v} j_e \vec{n}_e = 0$$



## Some issues about the Minkowski condition

- Minkowski condition:

$$\sum_{e \text{ at } v} j_e \vec{n}_e = 0$$

- The volume operator under this condition:

$$V|j_e \vec{n}_e\rangle \sim \sqrt{\epsilon^{ee'e''} j_e j_{e'} j_{e''} \vec{n}_e \cdot (\vec{n}_{e'} \times \vec{n}_{e''})} |j_e \vec{n}_e\rangle + o(\sqrt{j})$$

coincide with the classical expression.

## Some issues about the Minkowski condition

- Minkowski condition:

$$\sum_{e \text{ at } v} j_e \vec{n}_e = 0$$

- The volume operator under this condition:

$$V|j_e \vec{n}_e\rangle \sim \sqrt{\epsilon^{ee'e''} j_e j_{e'} j_{e''} \vec{n}_e \cdot (\vec{n}_{e'} \times \vec{n}_{e''})} |j_e \vec{n}_e\rangle + o(\sqrt{j})$$

coincide with the classical expression.

- The operator  $\sum_{ee' \text{ at } v} H_{v,ee'}$  doesn't preserve the condition.

## Some issues about the Minkowski condition

- Minkowski condition:

$$\sum_{e \text{ at } v} j_e \vec{n}_e = 0$$

- The volume operator under this condition:

$$V|j_e \vec{n}_e\rangle \sim \sqrt{\epsilon^{ee'e''} j_e j_{e'} j_{e''} \vec{n}_e \cdot (\vec{n}_{e'} \times \vec{n}_{e''})} |j_e \vec{n}_e\rangle + o(\sqrt{j})$$

coincide with the classical expression.

- The operator  $\sum_{ee' \text{ at } v} H_{v,ee'}$  doesn't preserve the condition.
- Does the evolution operator  $e^{iHt}$  preserve this condition?



## A toy model for LQG cosmology

- Regardless of the Lorentz part of the Hamiltonian operator,

## A toy model for LQG cosmology

- ▶ Regardless of the Lorentz part of the Hamiltonian operator,
- ▶ Consider the "continuous" limit of the above model

## A toy model for LQG cosmology

- ▶ Regardless of the Lorentz part of the Hamiltonian operator,
- ▶ Consider the "continuous" limit of the above model
- ▶ The Hilbert space  $\mathcal{H}_{\text{cos}} =$

$$\overline{\text{span}(\langle \vec{g} | \vec{j} \rangle := \bigotimes_{e \in \gamma} D_{j_e j_e}^{j_e} (n_e^{-1} g n_e))} \rightarrow L^2((\mathbb{R}^+)^N, d\vec{X})$$

## A toy model for LQG cosmology

- ▶ Regardless of the Lorentz part of the Hamiltonian operator,
- ▶ Consider the "continuous" limit of the above model
- ▶ The Hilbert space  $\mathcal{H}_{\text{cos}} =$   

$$\overline{\text{span}(\langle \vec{g} | \vec{j} \rangle := \bigotimes_{e \in \gamma} D_{j_e j_e}^{j_e} (n_e^{-1} g n_e))} \rightarrow L^2((\mathbb{R}^+)^N, d\vec{X})$$
- ▶ The Hamiltonian operator  $H = \sum_v \sqrt{\sum_{e,e'} H_{v,ee'}^E} \rightarrow H =$   

$$\sqrt{|\sum_{ee'} H_e^c H_{e'}^c|} = \sqrt{\sum_{ee'} \sqrt{|x_e x_{e'}|} \frac{\partial^2}{\partial x_e \partial x_{e'}} \sqrt{|x_e x_{e'}|}}$$

## Semiclassical analysis:

- Coherent state  $\Psi(\vec{x}) = \prod_{e \in \gamma} \psi(x_e)$

$$\psi(x) := \int_{-\infty}^{\infty} d\omega e^{-\frac{(\omega - \omega_0)^2}{2\sigma^2} + i\xi_0 \omega} \phi_{\omega}(x) = \sqrt{\frac{\sigma^2}{x}} e^{-\frac{\sigma^2}{2}(\xi_0 - \ln(x))^2 - i\omega_0(\xi_0 - \ln(x))}$$



## Semiclassical analysis:

- ▶ Coherent state  $\Psi(\vec{x}) = \prod_{e \in \gamma} \psi(x_e)$

$$\psi(x) := \int_{-\infty}^{\infty} d\omega e^{-\frac{(\omega - \omega_0)^2}{2\sigma^2} + i\xi_0 \omega} \phi_{\omega}(x) = \sqrt{\frac{\sigma^2}{x}} e^{-\frac{\sigma^2}{2}(\xi_0 - \ln(x))^2 - i\omega_0(\xi_0 - \ln(x))}$$

- ▶ The semiclassical condition:  $\xi_0 \gg 1$ ,  $\sigma \ll 1$ , while  $\sigma\xi_0 \gg 1$  for  $\Delta \ln x / \ln x \ll 1$ .

## Semiclassical analysis:

- ▶ Coherent state  $\Psi(\vec{x}) = \prod_{e \in \gamma} \psi(x_e)$

$$\psi(x) := \int_{-\infty}^{\infty} d\omega e^{-\frac{(\omega - \omega_0)^2}{2\sigma^2} + i\xi_0 \omega} \phi_{\omega}(x) = \sqrt{\frac{\sigma^2}{x}} e^{-\frac{\sigma^2}{2}(\xi_0 - \ln(x))^2 - i\omega_0(\xi_0 - \ln(x))}$$

- ▶ The semiclassical condition:  $\xi_0 \gg 1$ ,  $\sigma \ll 1$ , while  $\sigma\xi_0 \gg 1$  for  $\Delta \ln x / \ln x \ll 1$ .
- ▶ Solving the dynamic:

$$\Psi(\vec{x}, \tau) = e^{iH\tau} \Psi(\vec{x}) = \frac{1}{\sqrt{\prod_i x_i}} \int_{-\infty}^{\infty} d^N \omega e^{-\frac{\sum_i (\omega_i - \omega_0)^2}{2\sigma^2} + i \sum_i (\xi_0 - \ln(x_i)) \omega_i + i \sqrt{|\sum_{i \neq j} \omega_i \omega_j|} \tau}$$

## Semiclassical analysis:

- ▶ Coherent state  $\Psi(\vec{x}) = \prod_{e \in \gamma} \psi(x_e)$

$$\psi(x) := \int_{-\infty}^{\infty} d\omega e^{-\frac{(\omega - \omega_0)^2}{2\sigma^2} + i\xi_0 \omega} \phi_{\omega}(x) = \sqrt{\frac{\sigma^2}{x}} e^{-\frac{\sigma^2}{2}(\xi_0 - \ln(x))^2 - i\omega_0(\xi_0 - \ln(x))}$$

- ▶ The semiclassical condition:  $\xi_0 \gg 1$ ,  $\sigma \ll 1$ , while  $\sigma\xi_0 \gg 1$  for  $\Delta \ln x / \ln x \ll 1$ .
- ▶ Solving the dynamic:

$$\Psi(\vec{x}, \tau) = e^{iH\tau} \Psi(\vec{x}) = \frac{1}{\sqrt{\prod_i x_i}} \int_{-\infty}^{\infty} d^N \omega e^{-\frac{\sum_i (\omega_i - \omega_0)^2}{2\sigma^2} + i \sum_i (\xi_0 - \ln(x_i)) \omega_i + i \sqrt{|\sum_{i \neq j} \omega_i \omega_j|} \tau}$$

- ▶ The trajectory where the state peaks:

$$\ln x_i - \xi_0 = \frac{N-1}{2\sqrt{C_N^2}} \tau$$



## Conclusion and outlook

- ▶ Semiclassically, the quantum dynamic gives us an expanding universe.
- ▶ The peak satisfies the Minkowski condition.
- ▶ Future works: the quantum phenomenon, generalization to general graph, same problem with graph changing Hamiltonian.....



Thanks!